

## Derivation of Superelastic Cross Section

A “superelastic” collision (or collision of the second kind) occurs when an electron de-excites a higher lying state. In doing so, the electron gains energy

### Inelastic



### Superelastic



where  $\Delta\varepsilon$  is the energy loss (inelastic) or energy gain (superelastic) resulting from the collision. The superelastic cross section is obtained from the excitation cross section through the Klein-Rosseland formula.

Assume we have two atomic levels

$n_i$  = density of level

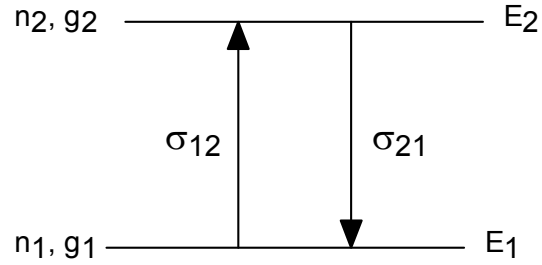
$g_i$  = degeneracy

$E_i$  = energy of level

$\Delta\varepsilon = E_2 - E_1$

$\sigma_{12}$  = excitation cross section

$\sigma_{21}$  = superelastic cross section



For a Maxwell Boltzmann electron energy distribution the number of electrons in  $(\varepsilon, \varepsilon + d\varepsilon)$  is proportional to

$$f(\varepsilon)d\varepsilon = K \exp\left(\frac{-\varepsilon}{kT_e}\right) \varepsilon^{1/2} d\varepsilon$$

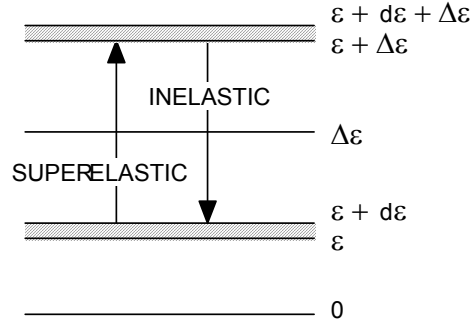
where  $K$  is a constant. The total rate of excitation of  $n_2$  by these electrons is

$$R_{12}(\varepsilon)d\varepsilon = K\sigma_{12}(\varepsilon)n_1\left(\frac{2\varepsilon}{m_e}\right)^{1/2} f(\varepsilon)d\varepsilon \quad \frac{\#}{\text{cm}^3\text{s}}$$

The total rate of de-excitation is

$$R_{21}(\varepsilon)d\varepsilon = K\sigma_{21}(\varepsilon)n_2\left(\frac{2\varepsilon}{m_e}\right)^{1/2} f(\varepsilon)d\varepsilon \quad \frac{\#}{\text{cm}^3\text{s}}$$

If there is only one inelastic process, then electrons can only enter  $(\varepsilon, \varepsilon + d\varepsilon)$  with  $\varepsilon < \Delta\varepsilon$  by suffering inelastic collisions from higher energies, and can only leave  $(\varepsilon, \varepsilon + d\varepsilon)$  by suffering superelastic collisions.



So in the steady state we must have the rate of population and depopulation of  $(\varepsilon, \varepsilon + d\varepsilon)$  be equal.

$$\phi_{\text{INTO}(\varepsilon, \varepsilon + d\varepsilon)} = R_{12}(\varepsilon + \Delta\varepsilon) \quad \left[ \begin{array}{l} \text{Electrons at } \varepsilon + \Delta\varepsilon \\ \text{exciting } n_1 \text{ and} \\ \text{losing energy.} \end{array} \right]$$

$$\phi_{\text{OUTOF}(\varepsilon, \varepsilon + d\varepsilon)} = R_{21}(\varepsilon) \quad \left[ \begin{array}{l} \text{Electrons at } \varepsilon \text{ de exciting } n_2 \\ \text{and gaining energy.} \end{array} \right]$$

$$\sigma_{12}(\varepsilon + \Delta\varepsilon) n_1 \left( \frac{2(\varepsilon + \Delta\varepsilon)}{m_e} \right)^{1/2} \exp\left( \frac{-(\varepsilon + \Delta\varepsilon)}{kT_e} \right) (\varepsilon + \Delta\varepsilon)^{1/2} =$$

$$\sigma_{21}(\varepsilon) n_2 \left( \frac{2\varepsilon}{m_e} \right)^{1/2} \exp\left( \frac{-\varepsilon}{kT_e} \right) \varepsilon^{1/2}$$

$$\sigma_{21}(\varepsilon) = \sigma_{12}(\varepsilon + \Delta\varepsilon) \left( \frac{\varepsilon + \Delta\varepsilon}{\varepsilon} \right) \frac{n_1}{n_2} \exp\left( \frac{-\Delta\varepsilon}{kT_e} \right)$$

If we are in equilibrium then  $\frac{n_2}{n_1} = \frac{g_2}{g_1} \exp\left( \frac{-\Delta\varepsilon}{kT_e} \right)$  so

$$\sigma_{21}(\varepsilon) = \frac{g_1}{g_2} \left( \frac{\varepsilon + \Delta\varepsilon}{\varepsilon} \right) \sigma_{12}(\varepsilon + \Delta\varepsilon)$$

$\sigma_{21}(\varepsilon)$  is always nonzero for  $\varepsilon > 0$  since  $\sigma_{12}(\varepsilon + \Delta\varepsilon)$  is nonzero.