MODELING OF MICRODISCHARGES FOR USE AS MICROTHRUSTERS

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AGENDA

- Introduction to microdischarge (MD) devices
- Description of model
- Reactor geometry and parameters
- Plasma characteristics
- Effect of geometry, and power
- Incremental thrust, and effect of power
- Concluding Remarks

MICRODISCHARGE PLASMA SOURCES

- Microdischarges are plasma devices which leverage pd scaling to operate dc atmospheric glows 10s $-100s \mu m$ in size.
- Few 100s V, a few mA
- Although similar to PDP cells, MDs are usually dc devices which largely rely on nonequilibrium beam components of the EED.
- Electrostatic nonequilibrium results from their small size. Debye lengths and cathode falls are commensurate with size of devices.



• Ref: Kurt Becker, GEC 2003

APPLICATIONS OF MICRODISCHARGES

- MEMS fabrication techniques enable innovative structures for displays and detectors.
- MDs can be used as microthrusters in small spacecraft for precise control which are requisites for array of satellites.







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Ref: http://www.design.caltech.edu/micropropulsion

DESCRIPTION OF MODEL

- To investigate microdischarge sources, nonPDPSIM, a 2dimensional plasma code was developed with added capabilities for pulsed operation.
 - Finite volume method in rectilinear or cylindrical unstructured meshes.
 - Implicit drift-diffusion-advection for charged species
 - Navier-Stokes for neutral species
 - Poisson's equation (volume, surface charge, material conduction)
 - Secondary electrons by impact, thermionics, photo-emission
 - Electron energy equation coupled with Boltzmann solution
 - Monte Carlo simulation for beam electrons.
 - Circuit, radiation transport and photoionization, surface chemistry models.

DESCRIPTION OF MODEL: CHARGED PARTICLE, SOURCES

 Continuity (sources from electron and heavy particle collisions, surface chemistry, photo-ionization, secondary emission), fluxes by modified Sharfetter-Gummel with advective flow field.

$$\frac{\partial N_i}{\partial t} = -\vec{\nabla} \cdot \vec{\phi} + S_i$$

- Poisson's Equation for Electric Potential: $-\nabla \cdot \mathcal{E} \nabla \Phi = \rho_V + \rho_S$
- Photoionization, electric field and secondary emission:

$$S_{Pi}(\vec{r}) = \int \frac{N_i(\vec{r})\sigma_{ij}N_j(\vec{r}')\exp\left(\frac{-\left|\vec{r}'-\vec{r}\right|}{\lambda}\right)d^3\vec{r}'}{4\pi\left|\vec{r}'-\vec{r}\right|^2}$$
$$S_{Si} = -\nabla \cdot j, \quad j_E = AT^2 \exp\left(\frac{-\left(\Phi_W - \left(q^3 E/\varepsilon_0\right)^{1/2}\right)}{kT_S}\right), \quad j_S = \sum_j \gamma_{ij}\phi_j$$

ELECTRON ENERGY, TRANSPORT COEFFICIENTS

• Bulk electrons: Electron energy equation with coefficients obtained from Boltzmann's equation solution for EED.

$$\frac{\partial (n_e \varepsilon)}{\partial t} = \vec{j} \cdot \vec{E} + \sigma E_{EM}^2 - n_e \sum_i N_i \kappa_i - \nabla \cdot \left(\frac{5}{2} \varepsilon \phi - \lambda \nabla T_e\right), \quad \vec{j} = q \vec{\phi}_e$$

$$\overset{\text{MCS region}}{\underset{\text{Plasma}}{\text{Cathode}}} \stackrel{\text{Plasma}}{\underset{\text{Cathode}}{\text{MCS mesh}}} \stackrel{\text{Cathode}}{\underset{\text{Cathode}}{\text{Cathode}}} \stackrel{\text{Beam Electrons: Monte Carlo Simulation}}{\underset{\text{Superimposed on unstructured fluid mesh. Construct Greens functions for interpolation between meshes.}}$$

DESCRIPTION OF MODEL: NEUTRAL PARTICLE TRANSPORT

• Fluid averaged values of mass density, mass momentum and thermal energy density obtained using unsteady, compressible algorithms.

$$\frac{\partial \rho}{\partial t} = -\nabla \cdot (\rho \vec{v}) + (inlets, pumps)$$

$$\frac{\partial (\rho \vec{v})}{\partial t} = \nabla \left(\sum_{i} N_{i} k T_{i} \right) - \nabla \cdot (\rho \vec{v} \vec{v}) - \vec{\nabla} \cdot \vec{\tau} + \sum_{i} \left(q_{i} N_{i} \vec{E}_{i} - S_{i} m_{i} \mu_{i} q_{i} \vec{E} \right)$$

$$\frac{\partial (\rho c_{p} T)}{\partial t} = -\nabla \left(-\kappa \nabla T + \rho \vec{v} c_{p} T \right) + P_{i} \nabla \cdot v_{f} - \sum_{i} R_{i} \Delta H_{i} + \sum_{i} \vec{j}_{i} \cdot \vec{E}$$

• Individual species are addressed with superimposed diffusive transport.

$$N_i(t + \Delta t) = N_i(t) - \nabla \cdot \left(\vec{v}_f - D_i N_T \nabla \left(\frac{N_i(t + \Delta t)}{N_T}\right)\right) + S_V + S_S$$

GEOMETRY AND MESH

- Plasma dia: 150 μm at inlet, 250 μm at cathode.
- \bullet Electrodes 130 μm thick.
- Dielectric gap 1.5 mm.
- Geometry B: 1.5 mm dielectric above the cathode.
- Fine meshing near electrodes, less refined near exit.
- Anode grounded; cathode bias varied based on power deposition (0.25 - 1.0 W).
- 10 sccm Ar, 30 Torr at inlet, 10 Torr at exit.



EXPERIMENT: GEOMETRY





- Modeled geometry similar to experimental setup.
- Plume characterized by densities of excited states.
 - Ref: John Slough, J.J. Ewing, AIAA 2005-4074

CHARGED SPECIES: GEOMETRY A



Power deposition occurs in the cathode fall by collisions with hot electrons.

• 10 sccm Ar, 0.5 W

• Very high electric fields near cathode.

NEUTRAL FLUID PROPERTIES: GEOMETRY A



- Plume extends downstream, can be used for diagnosis.
- Gas heating and consequent expansion is 10 sccm Ar, 30 10 Torr a source of thrust. • 0.5 W.
 - Ref: John Slough, J.J. Ewing, AIAA 2005-4074

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VELOCITY INCREASE WITH DISCHARGE



- Gas heating and subsequent expansion causes increase in velocity.
- Steady state after one or two bursts of flow.
- At high plasma density, momentum transfer between charged species and neutrals is also important.

- 10 sccm Ar, 30 Torr at inlet, 10 Torr at exit.
- 0.5 Watts.
- Power turned on at 0.5 ms.

POWER DEPOSITION: IONIZATION SOURCES



Bulk ionization (cm⁻³ sec⁻¹) 1 100 Beam ionization (cm⁻³ sec⁻¹)

- Ionization rates increase with power.
- Beam electrons are equally as important as bulk electrons.
- 10 sccm Ar, 30 Torr at inlet, 10 Torr at exit.

POWER DEPOSITION: PLASMA PROPERTIES



- Hotter gases lead to higher ΔV and higher thrust production.
- Increase in mean free path due to rarefaction may affect power deposited to neutrals.
- With increasing [e], increase in production of electronically excited states.
 - 10 sccm Ar, 30 Torr at inlet, 10 Torr at exit.

POWER DEPOSITION: FLOW VELOCITY



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BASE CASE RESULTS: GEOMETRY B



- Electrons are confined, discharge operates in an unsteady regime.
- Ionization pulses travel towards anode.
- Power densities are greater than that of Geometry A.
 - 10 sccm Ar, 30 10 Torr
 - 0.5 W, turned on at 0.5 ms

VELOCITY INCREASE: GEOMETRY B



POWER DEPOSITION: GEOMETRY B



- Discharge operates in normal glow, current increases with power, whereas voltage marginally increases.
- [e] increases substantially with increase in power.
- With increasing [e], charge buildup on the dielectric can be high.
- 10 sccm Ar, 30 Torr at inlet, 10 Torr at exit.

CURRENT VOLTAGE CHARACTERISTICS

- Operating voltage for geometry A remains almost a constant (260 V), whereas slight changes observed for geometry B.
- Discharge resistance R_D of 43 kΩ.



• Thrust calculated by:

$$F = \frac{dm}{dt}V_e + A_e(P_e - P_a)$$

• Increase in thrust is the rate of momentum transfer to the neutrals when the discharge is switched on.

$$\overrightarrow{\delta F} = \left(\frac{dm}{dt}\overrightarrow{V}\right)_{pulse} - \left(\frac{dm}{dt}\overrightarrow{V}\right)_{nopulse}$$

• Meaningful incremental thrust occurs when power deposited to plasma is greater than that contained in the flow.

$$\vec{j}.\vec{E} \ge \frac{1}{2}\dot{\rho}v^2$$

INCREMENTAL THRUST: EFFECT OF POWER



- 10 sccm Ar, 30 Torr upstream, 10 Torr downstream.
- Power turned on at 0.5 ms

CONCLUDING REMARKS

- An axially symmetric microdischarge was computationally investigated with potential application to microthrusters.
- Studies were conducted to investigate the effect of parameters such as power deposition, and the geometry of the reactor.
- The geometry affected the plasma characteristics significantly, whereas there was no significant difference to incremental thrust.
- At higher power, higher gas temperatures lead to higher thrust.
- Rarefaction at high temperatures decreases mean free path and could limit thrust produced.