

SCALING OF MICRODISCHARGE DEVICES: PYRAMIDAL STRUCTURES*

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AGENDA

- Introduction of pyramidal microdischarge devices.
- Description of model.
- Fundamental properties of MDs sustained in neon.
- Transition from Townsend to negative glow.
- Scaling of MDs
- Concluding remarks.

MICRODISCHARGE PLASMA SOURCES

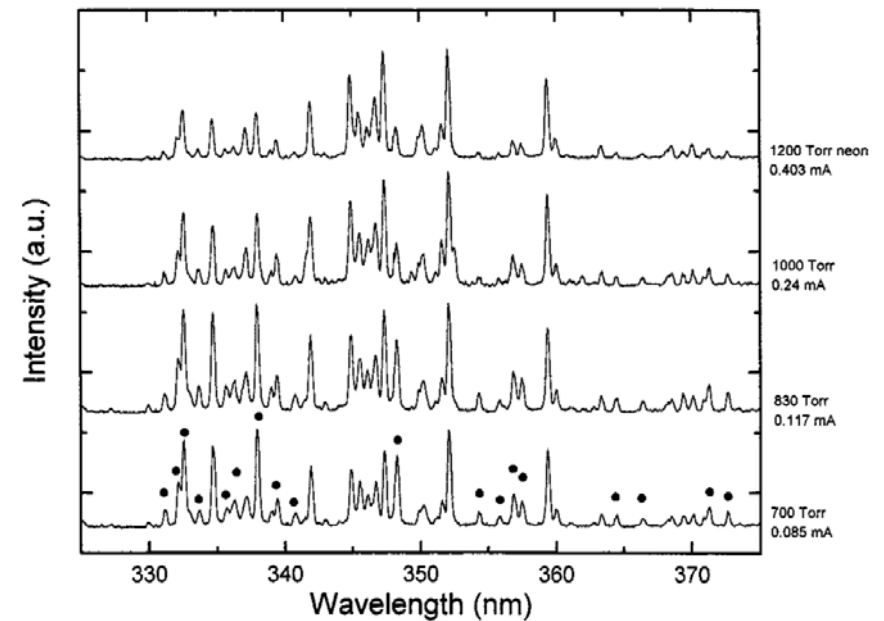
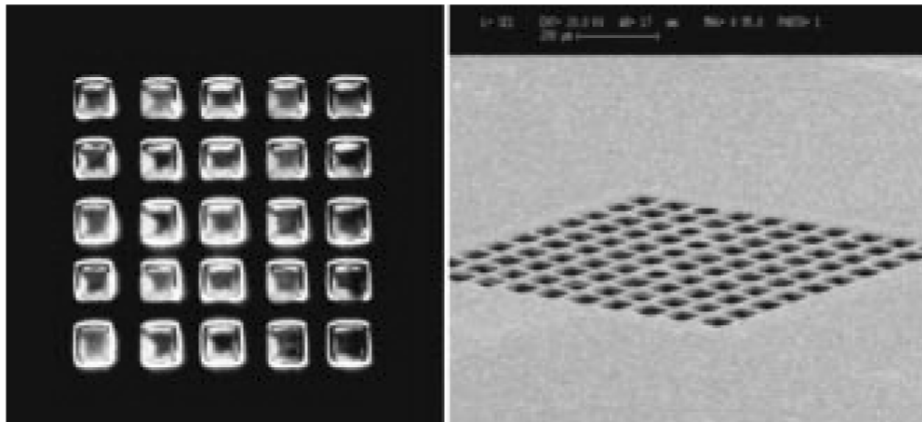
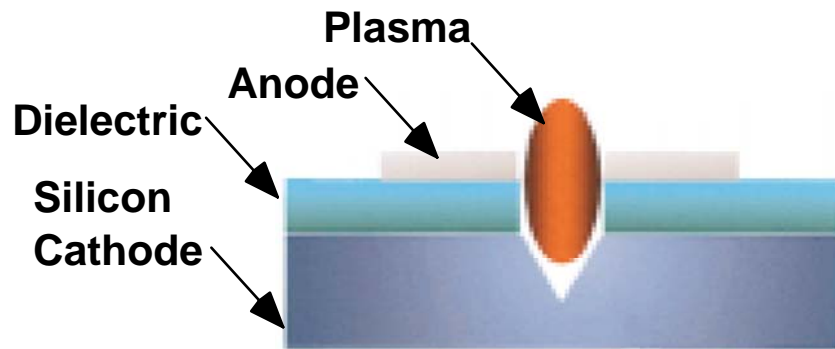
- Microdischarges are plasma devices which leverage pd scaling to operate dc atmospheric glows 10s –100s μm in size.
- MEMS fabrication techniques enable innovative structures for displays and detectors.
- Although similar to PDP cells, MDs are dc devices which largely rely on nonequilibrium beam components of the EED.
- Electrostatic nonequilibrium results from their small size. Debye lengths and cathode falls are commensurate with size of devices.

$$L_{\text{cathode Fall}} = \left(2V_c \epsilon_0 / (qn_I)\right)^{1/2} \approx 10 - 20 \mu\text{m}$$

$$\lambda_D \approx 750 \left(\frac{T_{eV}}{n_e (cm^{-3})} \right)^{1/2} cm \approx 10 \mu\text{m},$$

PYRAMIDAL MICRODISCHARGE DEVICES

- Si MDs with 10s μm pyramidal cavities display nonequilibrium behavior: Townsend to negative glow transitions.
- Small size also implies electrostatic nonequilibrium.



- S.-J. Park, et al., J. Sel. Topics Quant. Electron 8, 387 (2002); Appl. Phys. Lett. 78, 419 (2001).

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2-D MODELING OF MICRODISCHARGE SOURCES

- **Charged particle continuity (fluxes by Sharfetter-Gummel form)**

$$\frac{\partial N_i}{\partial t} = -\vec{\nabla} \cdot \left(qN_i \mu_i (-\vec{\nabla} \Phi) - D_i \nabla N_i \right) + S_i$$

- **Poisson's Equation for Electric Potential**

$$-\nabla \cdot \epsilon \nabla \Phi = \rho_V + \rho_S$$

- **Bulk continuum electron energy transport and MCS beam.**

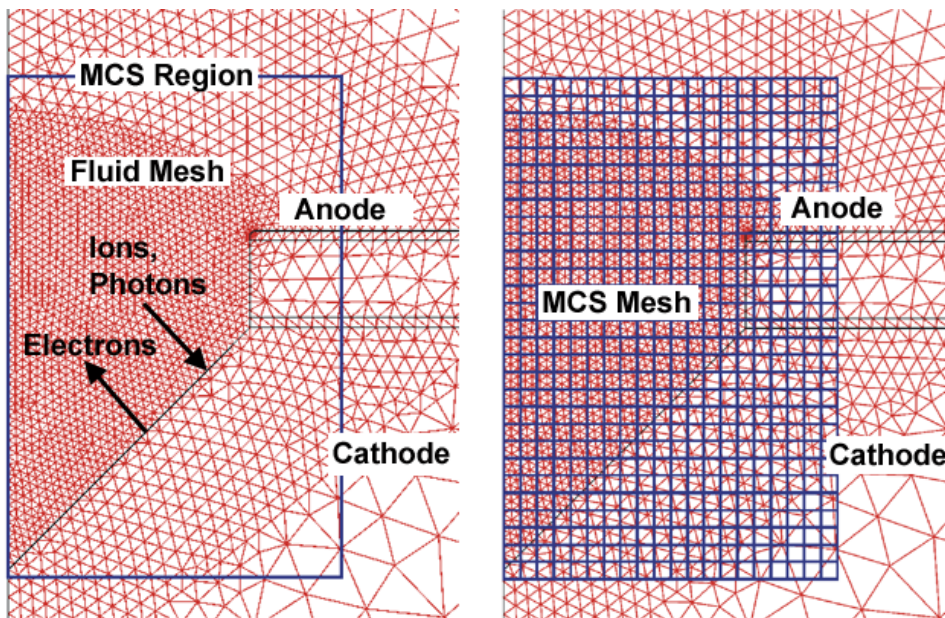
$$\frac{\partial (n_e \epsilon)}{\partial t} = \vec{j} \cdot \vec{E} - n_e \sum_i N_i \kappa_i - \nabla \cdot \left(\frac{5}{2} \epsilon \nabla \Phi - \lambda \nabla T_e \right), \quad \vec{j} = q \vec{\phi}_e$$

- **Neutral continuity and energy transport.**

$$\frac{\partial N_i}{\partial t} = -\nabla \cdot \left(\vec{v} - DN_o \nabla \left(\frac{N_i}{N_o} \right) \right) + S_i, \quad \frac{\partial (\rho c T)}{\partial t} = -\nabla \cdot \kappa \nabla T + P_g$$

DESCRIPTION OF MODEL: MCS AND MESHING

- Transport of energetic secondary electrons is addressed with a Monte Carlo Simulation.

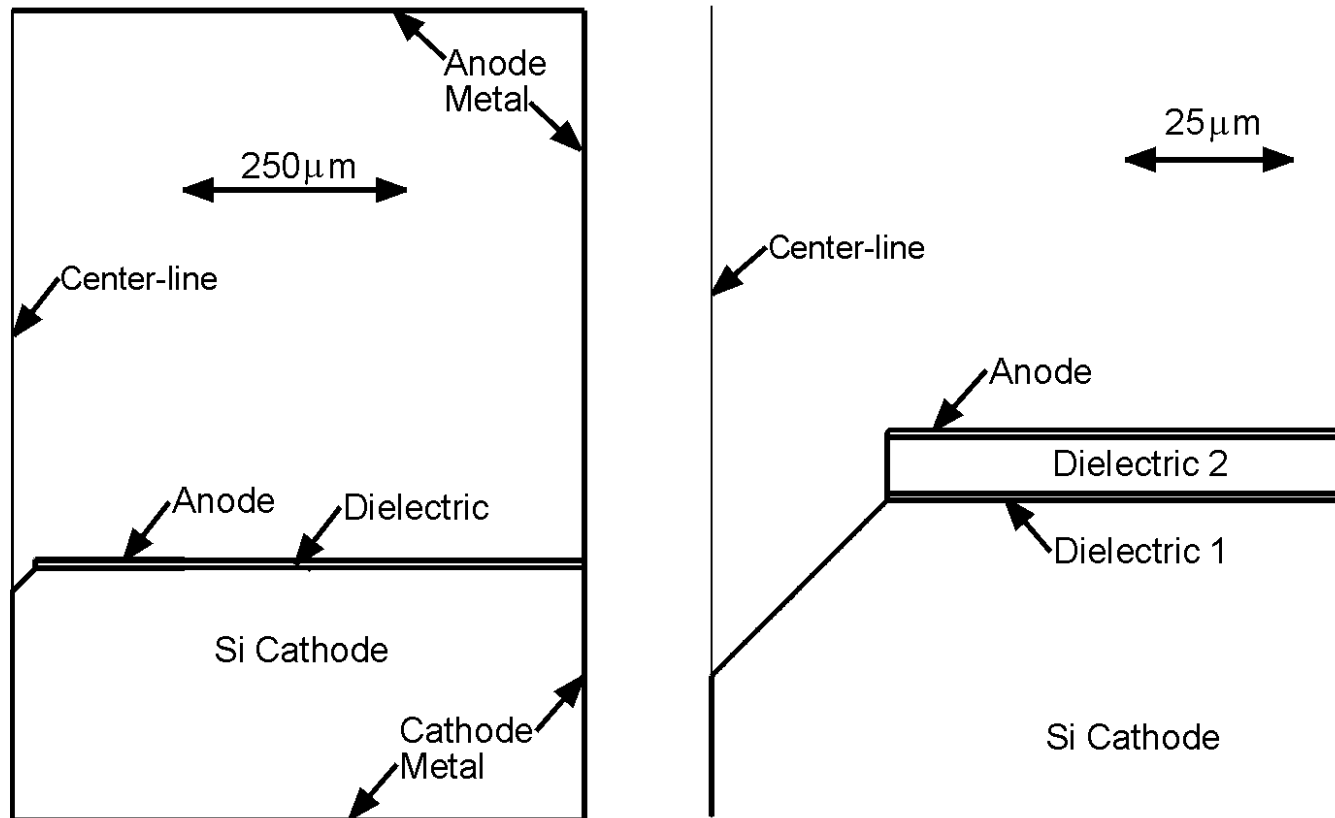


- Superimpose Cartesian MCS mesh on unstructured fluid mesh. Construct Greens functions for interpolation between meshes.
- Electrons and their progeny are followed until slowing into bulk plasma or leaving MCS volume.
- Electron energy distribution is computed on MCS mesh.
- EED produces source functions for electron impact processes which are interpolated to fluid mesh.

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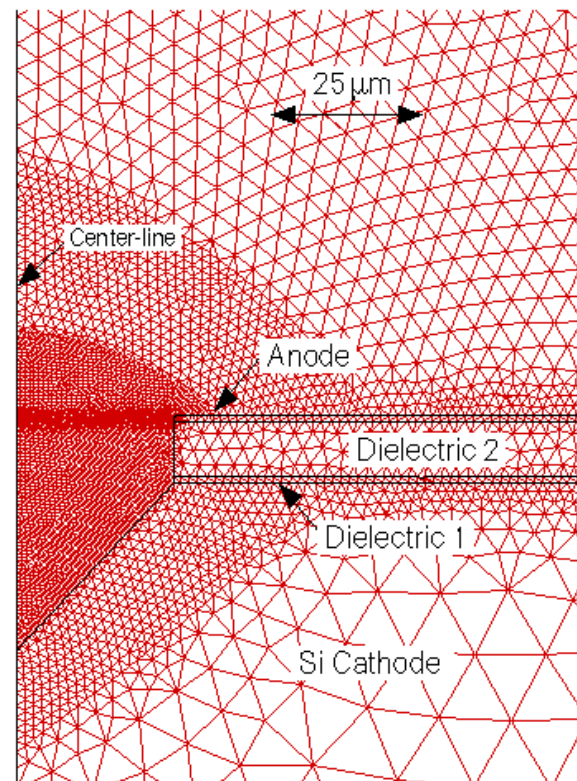
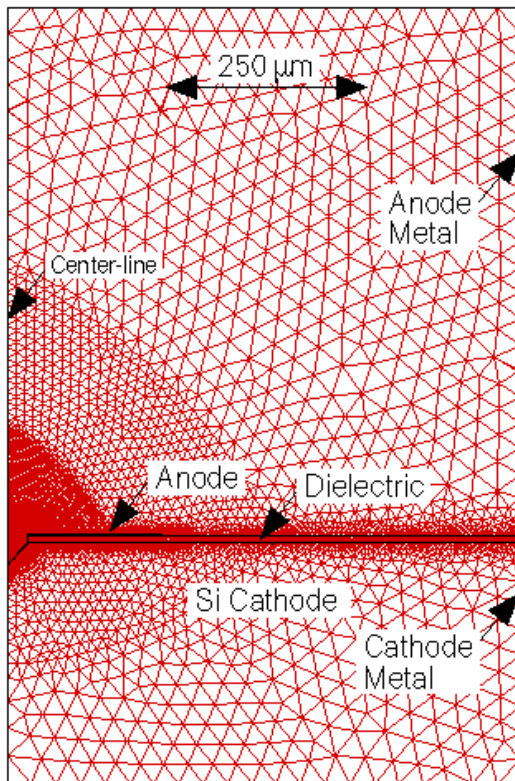
MODEL GEOMETRY: Si PYRAMID MICRODISCHARGE

- Investigations of a cylindrically symmetric Si pyramid microdischarge were performed.



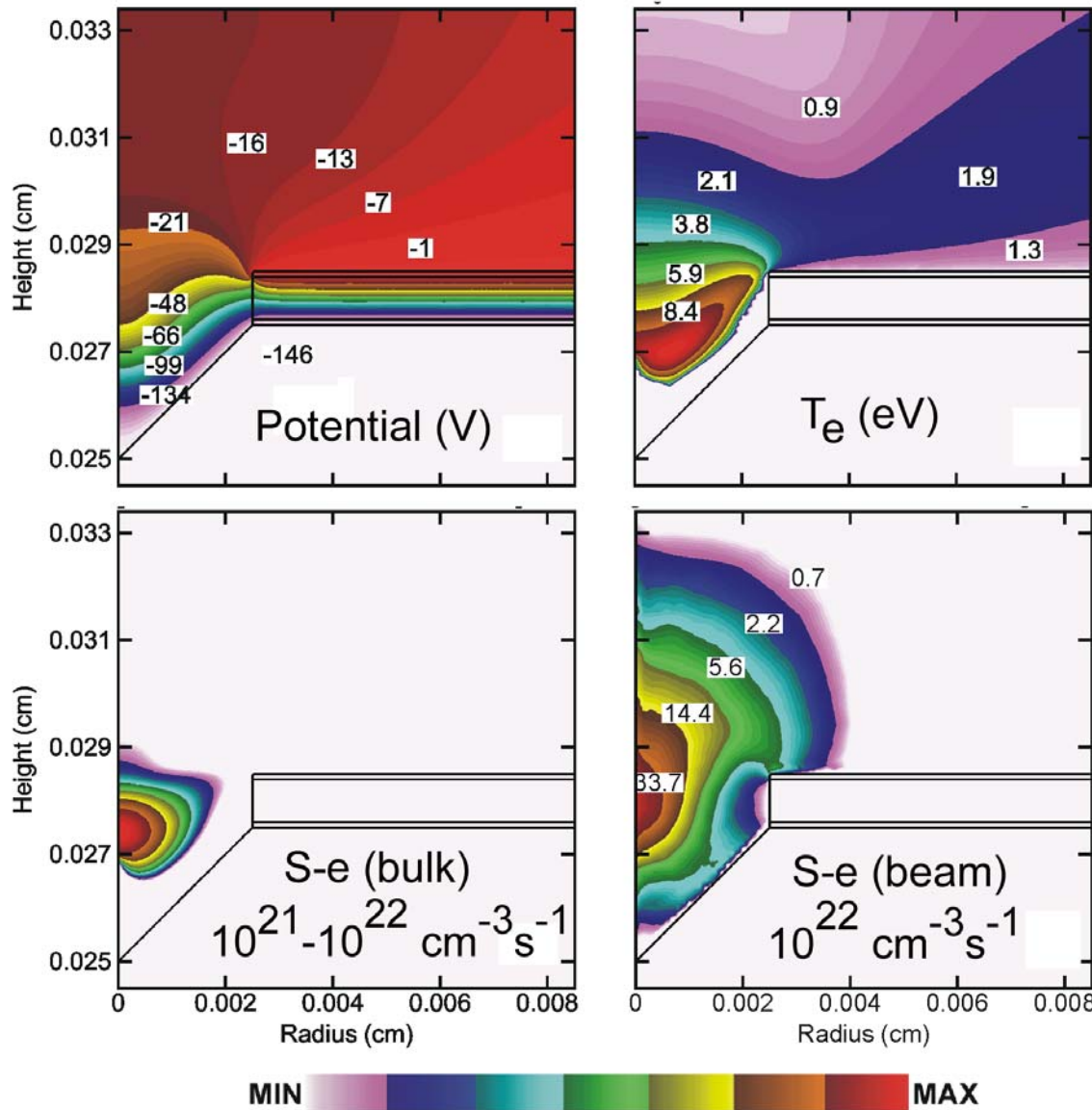
MODEL GEOMETRY: Si PYRAMID MICRODISCHARGE

- Meshing is absolutely critical to resolve small structures and distant boundaries.
- Typical Mesh: 5,000-10,000 nodes, 10^2 - 10^3 dynamic range



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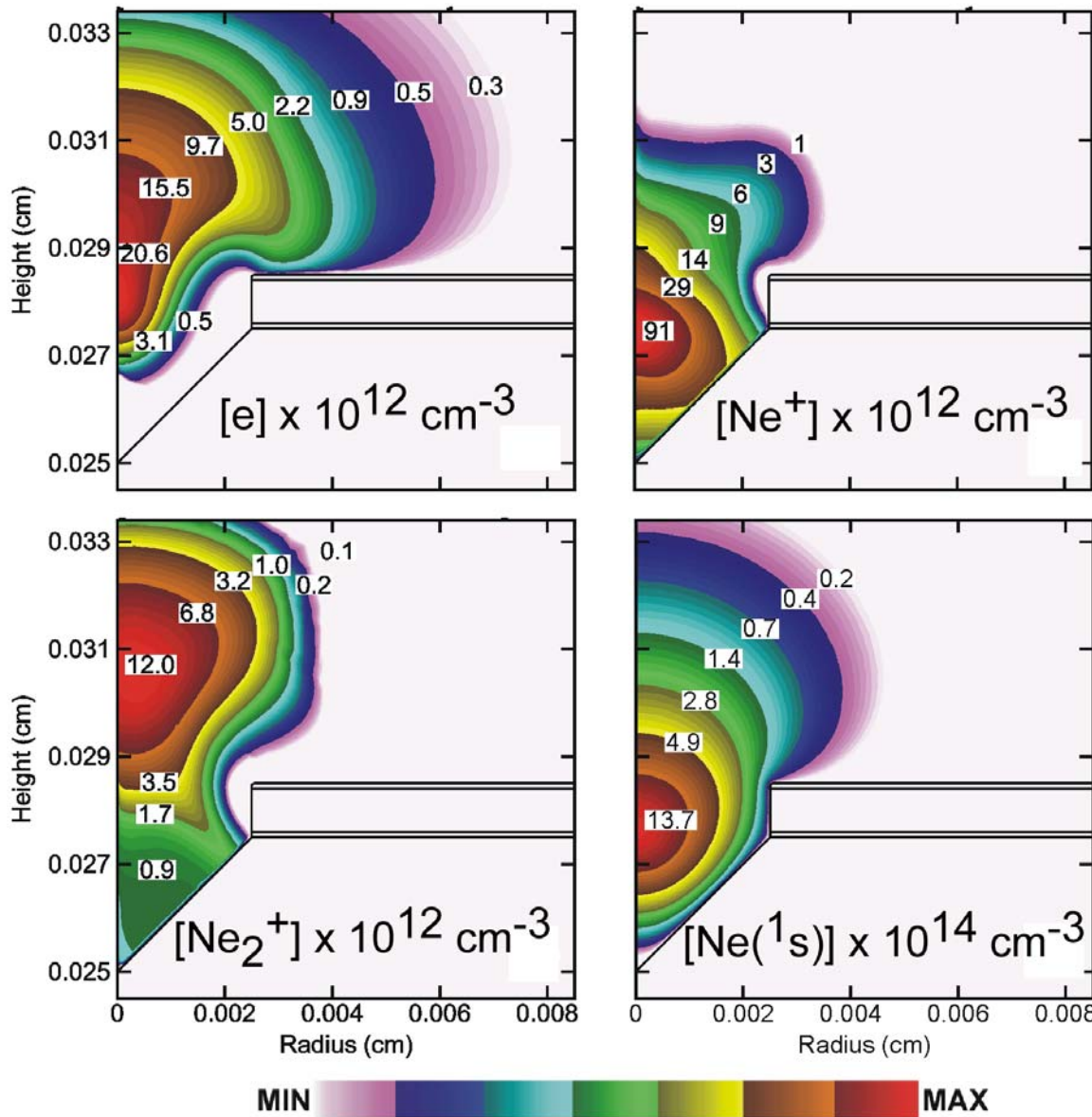
BASE CASE: Ne, 600 Torr, 50 μm DIAMETER



- Optimum operation produces large enough charge density to warp electric potential into cathode well.
- In spite of large T_e , ionization is dominated by beam electrons
- Ne, 600 Torr, 50 μm , 200 V, 1 $\text{M}\Omega$

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BASE CASE: Ne, 600 Torr 50 μm DIAMETER



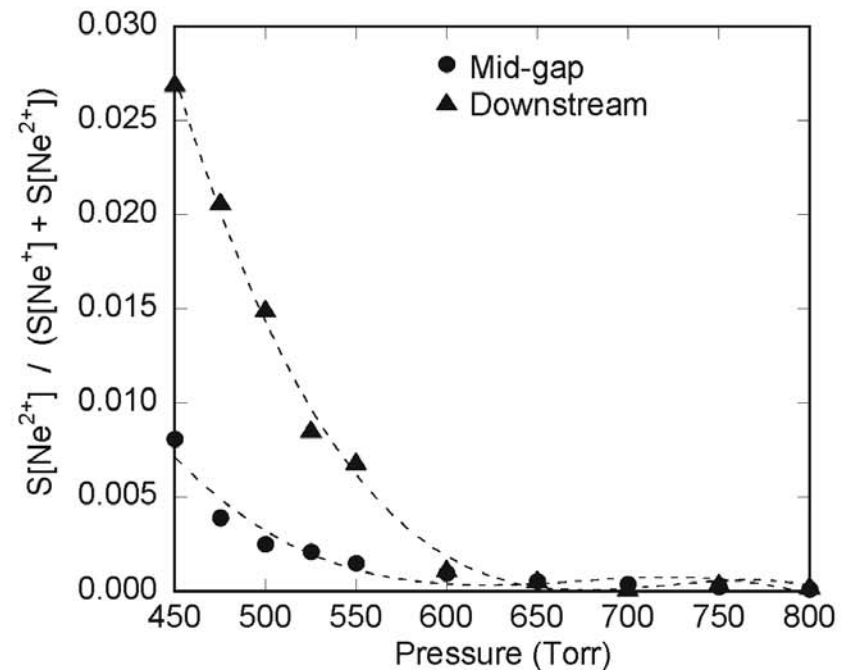
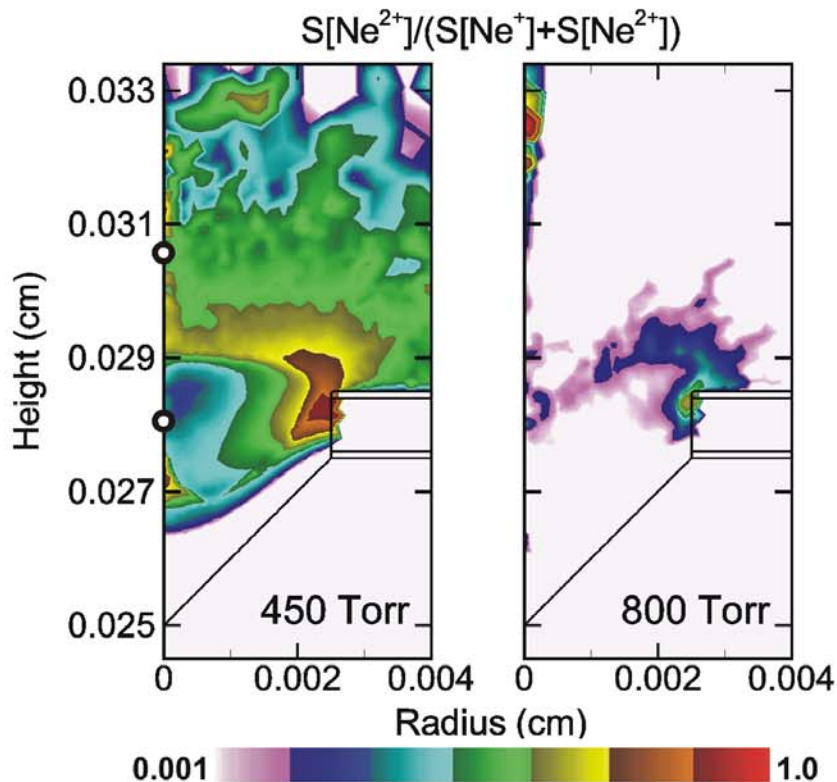
- There is essentially no region of quasi-neutrality or which is positive column-like.
- Monomer and dimer ions are segregated.
- Excited state densities $> 10^{15} \text{ cm}^{-3}$ rival macroscopic devices

- Ne, 600 Torr, 50 μm , 200 V, 1 $\text{M}\Omega$

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TRANSITION TO NEGATIVE-GLOW BEHAVIOR

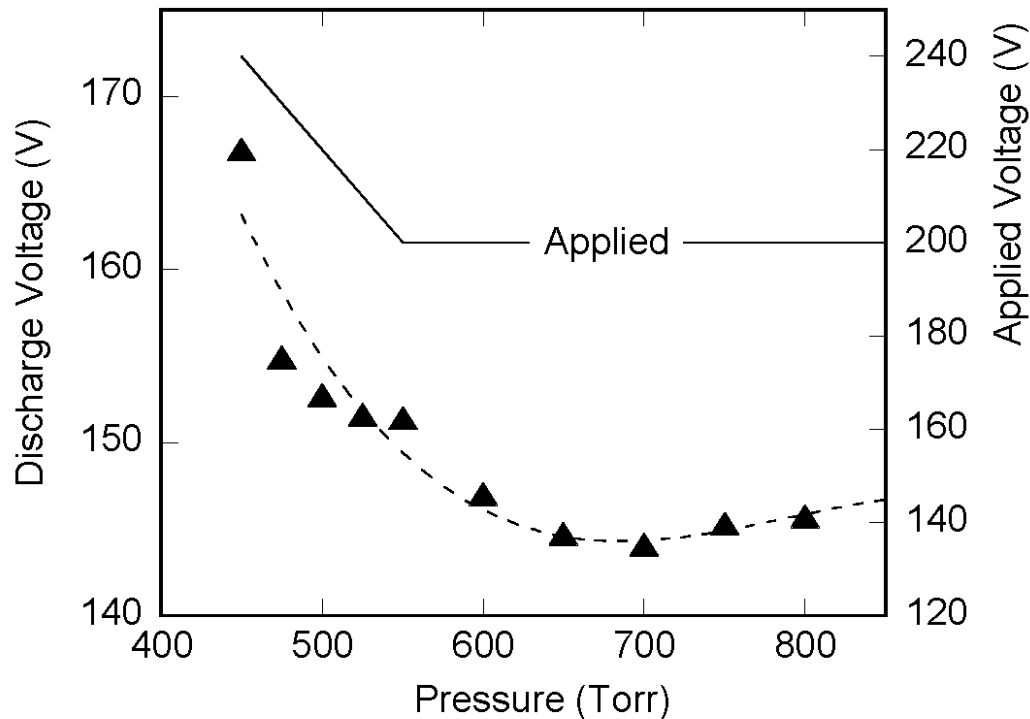
- Although geometry precludes true hollow cathode behavior, negative glow behavior sets in at lower pressures.
- Characterize negative glow by $S[\text{Ne}^{2+}] / (S[\text{Ne}^+] + S[\text{Ne}^{2+}])$



- Ne, 50 μm diameter, 200 V, 1 M Ω

SCALING WITH PRESSURE: PASCHEN BEHAVIOR

- With $pd=1-10$ Torr-cm, these microdischarge devices display Paschen behavior.



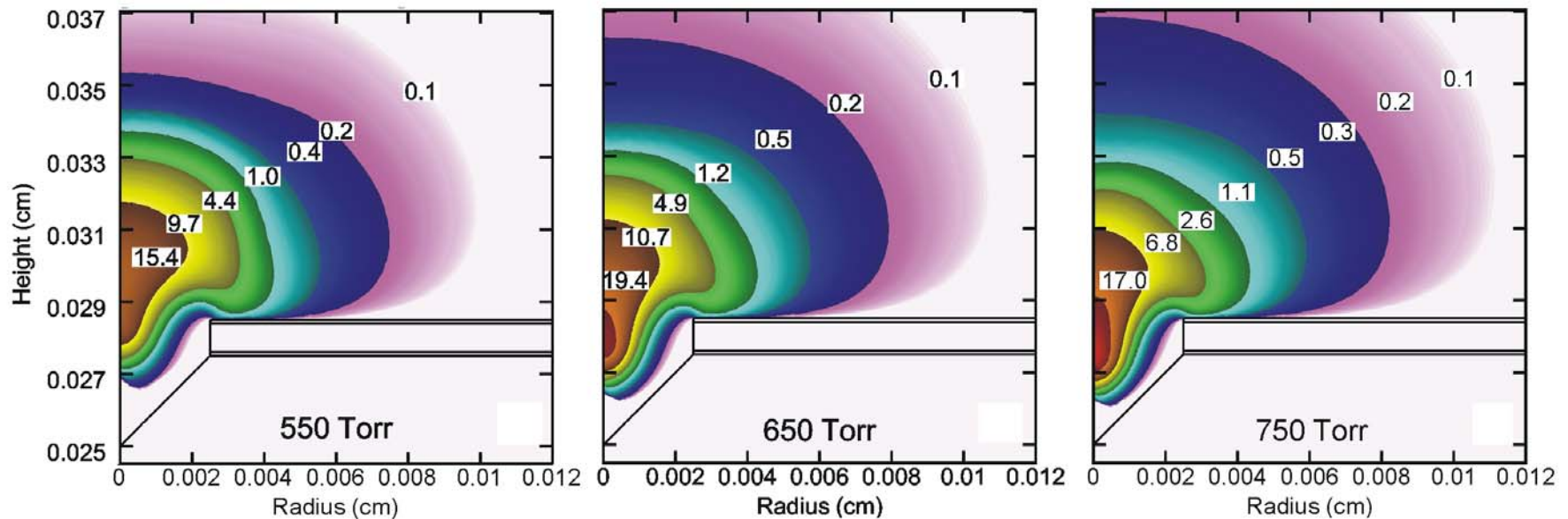
- Although sensitive to ballasting and current density lower pressures requiring larger applied voltages also produce large plasma densities.

- Ne, 50 μm diameter, 1 M Ω

SCALING WITH PRESSURE: PLASMA PROPERTIES

- Over a range of pressures that V (applied) and R (ballast) can be constant, confinement at higher pressures produces higher peak plasma densities.

• $[e] \times 10^{12} \text{ cm}^{-3}$



• **550 Torr**
 $[2.1 \times 10^{13} \text{ cm}^{-3}]$

• **650 Torr**
 $[3.9 \times 10^{13} \text{ cm}^{-3}]$

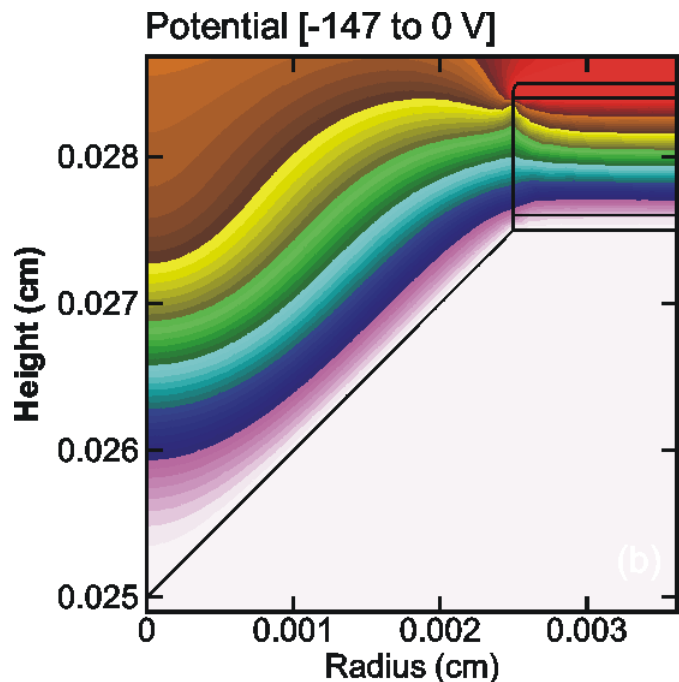
• **750 Torr**
 $[5.6 \times 10^{13} \text{ cm}^{-3}]$

• Ne, 50 μm diameter, 200V, 1 M Ω

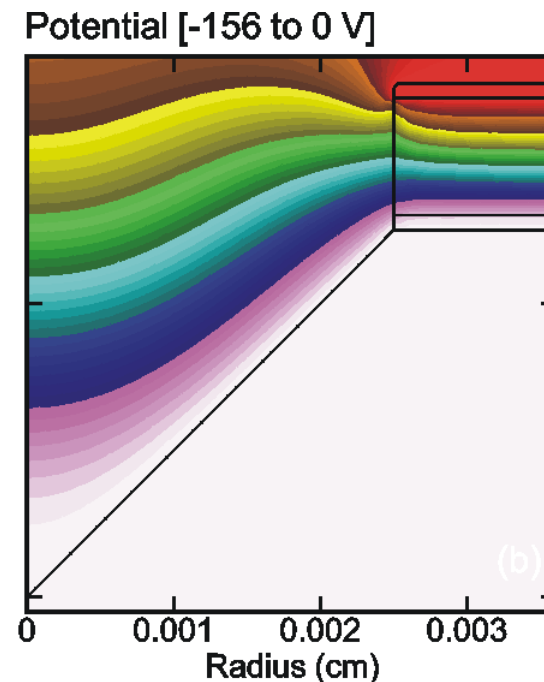
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SCALING CONSIDERATIONS: CATHODE FALL THICKNESS

- In MDs, the cathode fall thickness may be commensurate with cavity size. Current density is therefore critical to scaling.



- -210 V, 1 M Ω
[e]= $4.9 \times 10^{13} \text{ cm}^{-3}$



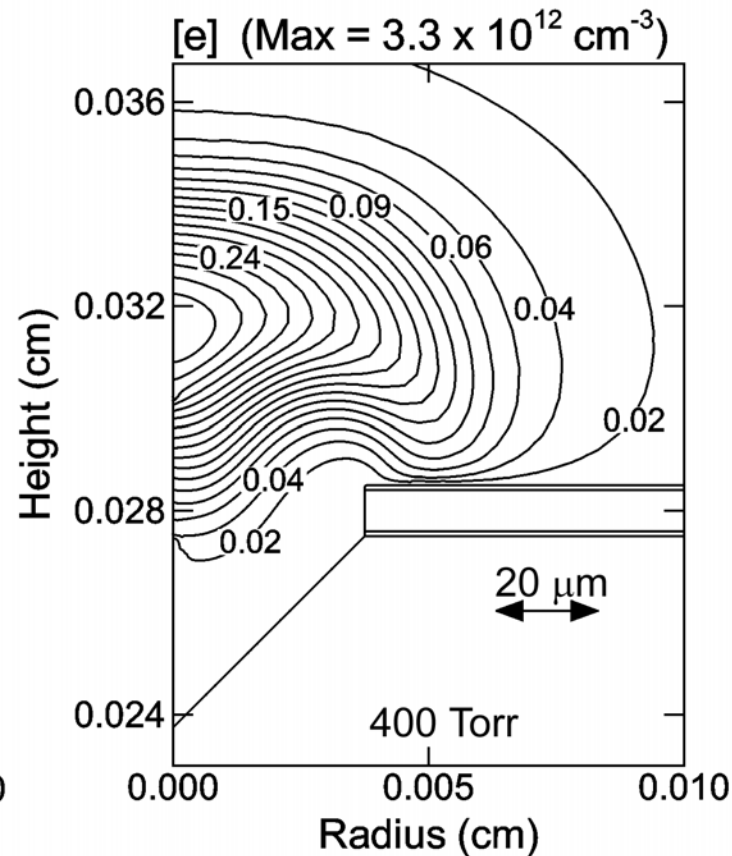
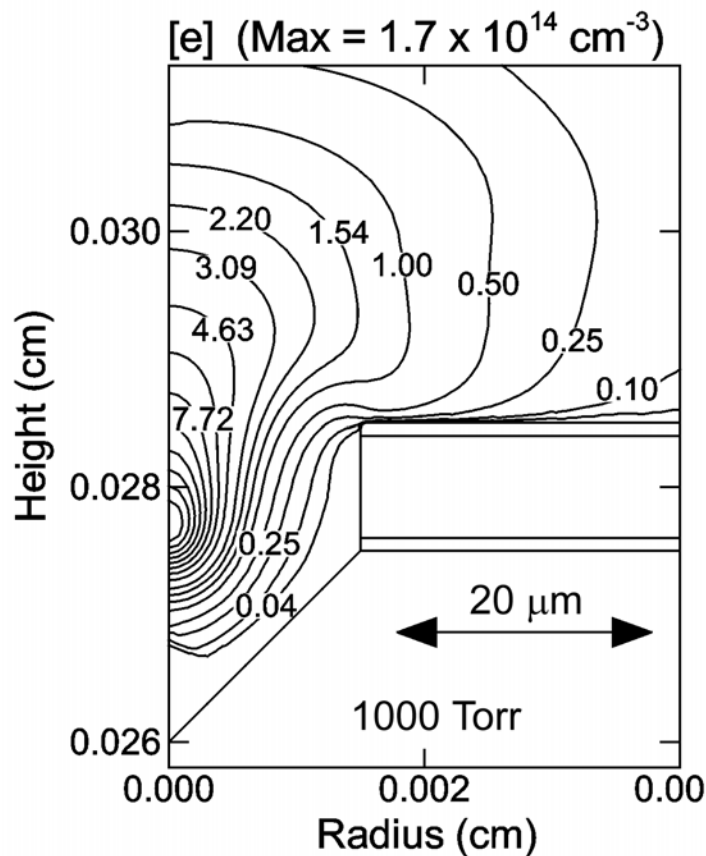
- -200 V, 1.75 M Ω
[e]= $5.3 \times 10^{12} \text{ cm}^{-3}$

- Low j (and [e]) may result in cathode fall not being conformal to cathode.

- Ne, 50 μm diameter, 600 Torr

SCALING WITH SIZE: pd, BALLAST = CONSTANT

- Scaling while maintaining pd, V(applied) and R(ballast) constant results in a reduced j and [e] in the larger device. The plasma is not conformal to the cathode.

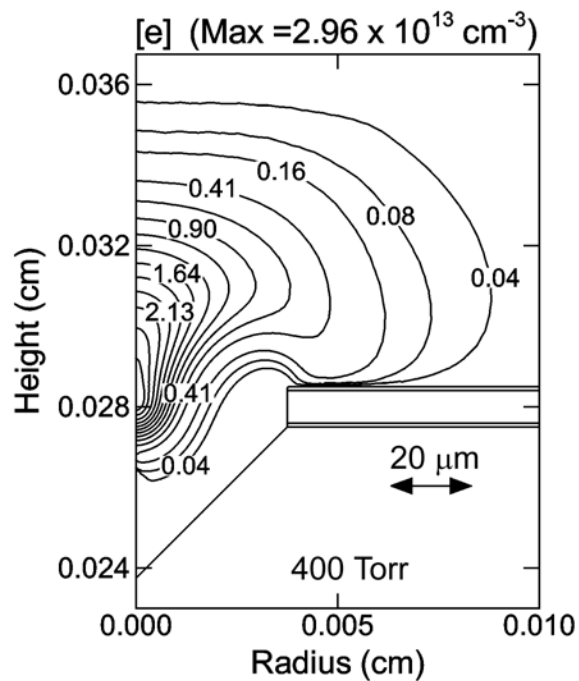


• Ne, -200 V, 1 M Ω

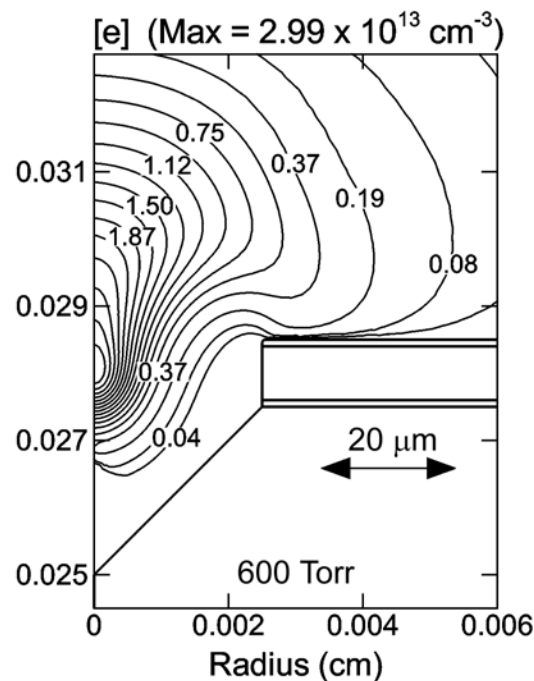
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SCALING WITH SIZE: pd , $j = \text{CONSTANT}$

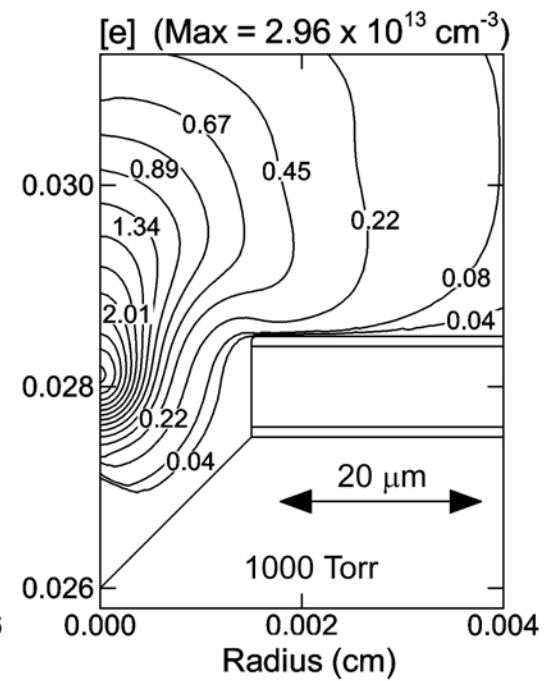
- Scaling while maintaining pd and j constant produces similar plasma densities and conformality to the cathode.



• 400 Torr



• 600 Torr



• 1000 Torr

• Ne, -200 V

CONCLUDING REMARKS

- MDs differ from macroscopic devices in that plasma scale lengths are commensurate with device dimensions.
- Scaling of MDs with pressure (traditionally “pd”) likely also required λ/L to remain constant or less than a critical value.
- Scaling with complex shapes must consider all dimensions.
- The transition from Townsend to negative glow is largely geometrically dependent, and can be controlled to some degree by shape.