

CONSEQUENCES OF RADIATION TRAPPING ON ELECTRON ENERGY DISTRIBUTIONS IN LOW PRESSURE INDUCTIVELY COUPLED Hg/Ar DISCHARGES*

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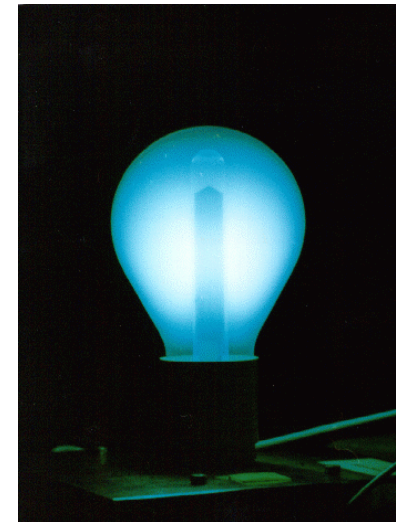
*** Work supported by Osram Sylvania Inc., EPRI and the NSF**

AGENDA

- **Radiation transport in low pressure plasmas**
- **Overview of the Hybrid Plasma Equipment Model**
- **The Monte Carlo Radiation Transport Model**
- **Effect of radiation trapping on EEDFs**
- **Variation of EEDF with**
 - **Power**
 - **Frequency**
 - **Radiative lifetime**
- **Conclusions**

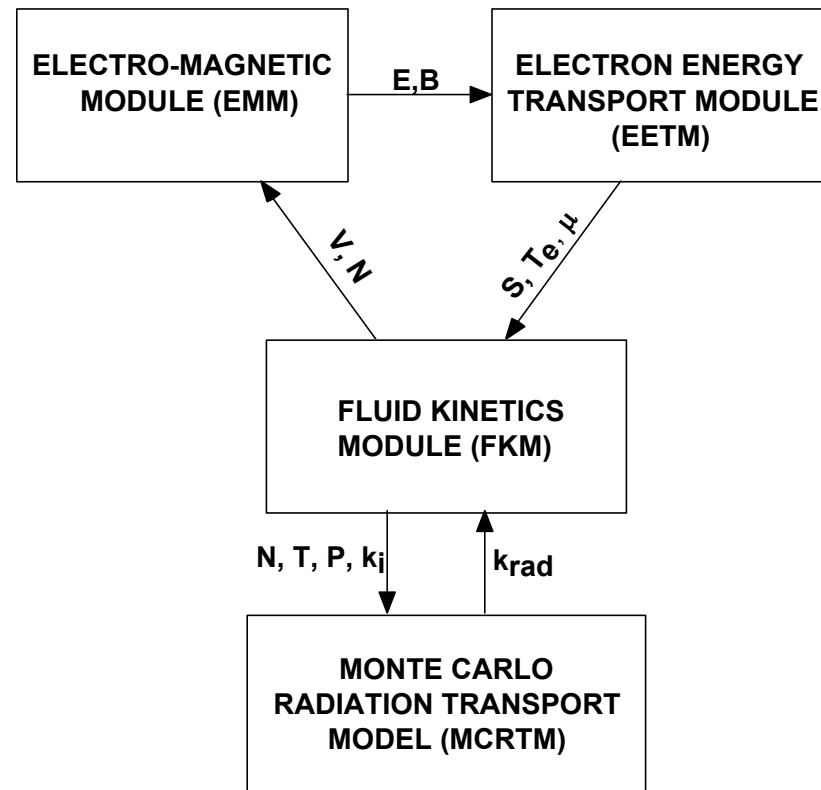
ELECTRODELESS LAMPS AND TRAPPING

- Electrodeless gas discharges are finding increasing use in the lighting industry due to their increased lifetime.
- Investigations are underway to increase the efficiency of these lamps, now $\cong 25\%$.
- Typical fluorescent lamps consist of Ar/Hg $\approx 97/3$. Resonance radiation from Hg (6^3P_1) (254 nm) and Hg (6^1P_1) (185 nm) excites phosphors which generate visible light.
- This resonance radiation can be absorbed and reemitted many times prior to striking the phosphor, increasing the effective lifetime of emission as viewed from outside the lamp.
- We have modeled this mechanism using a Monte Carlo radiation transport model linked to a hybrid plasma equipment model, to realistically simulate the discharge.

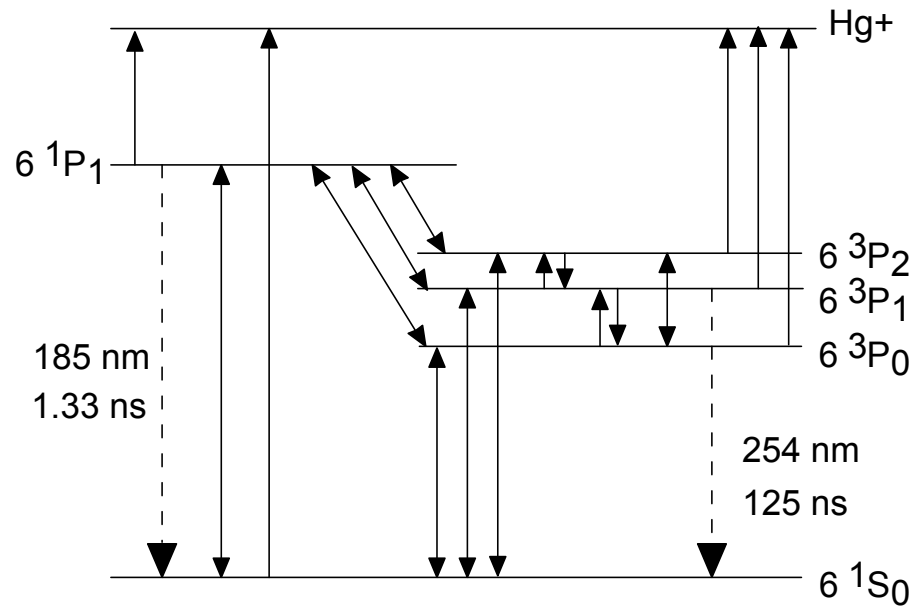
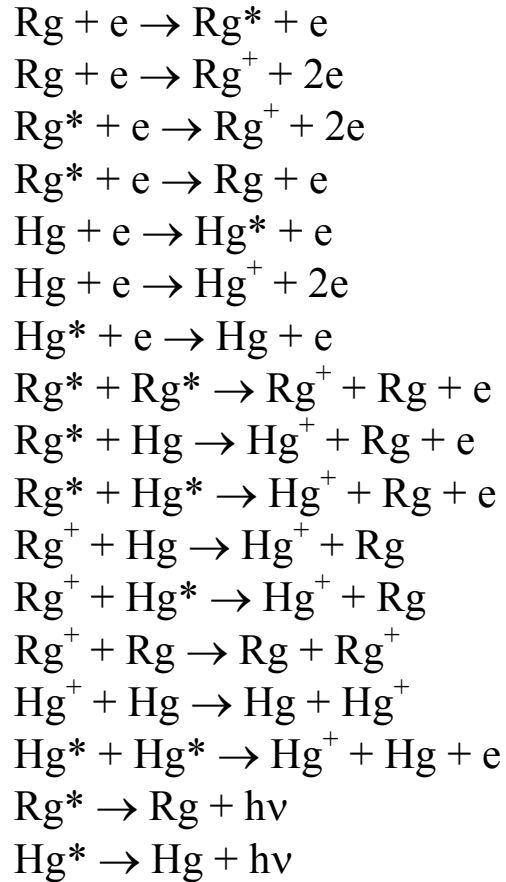


HYBRID PLASMA EQUIPMENT MODEL

- A modular simulator addressing low temperature, low pressure plasmas.
- **EMM**: electromagnetic fields and magneto-static fields
- **EETM**: electron temperature, electron impact sources, and transport coefficients
- **FKM**: densities, momenta, and temperatures of charged and neutral plasma species; and electrostatic potentials



REACTION CHEMISTRY

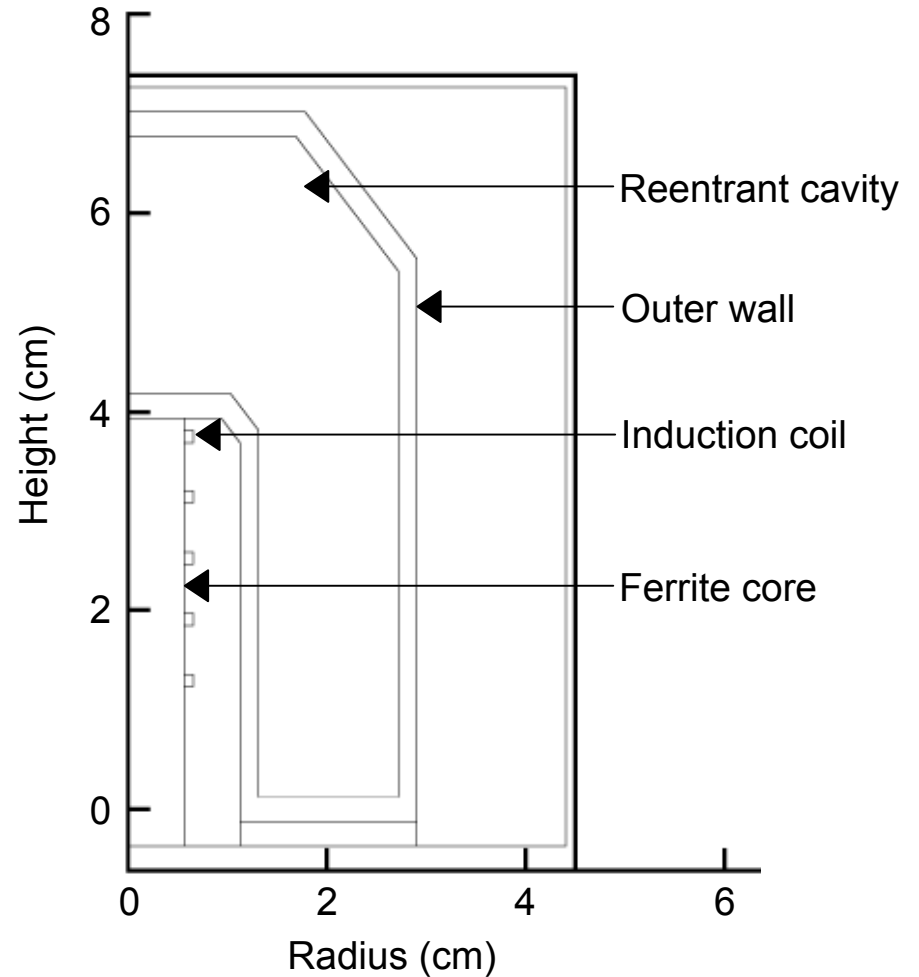


MONTE CARLO RADIATION TRANSPORT MODULE

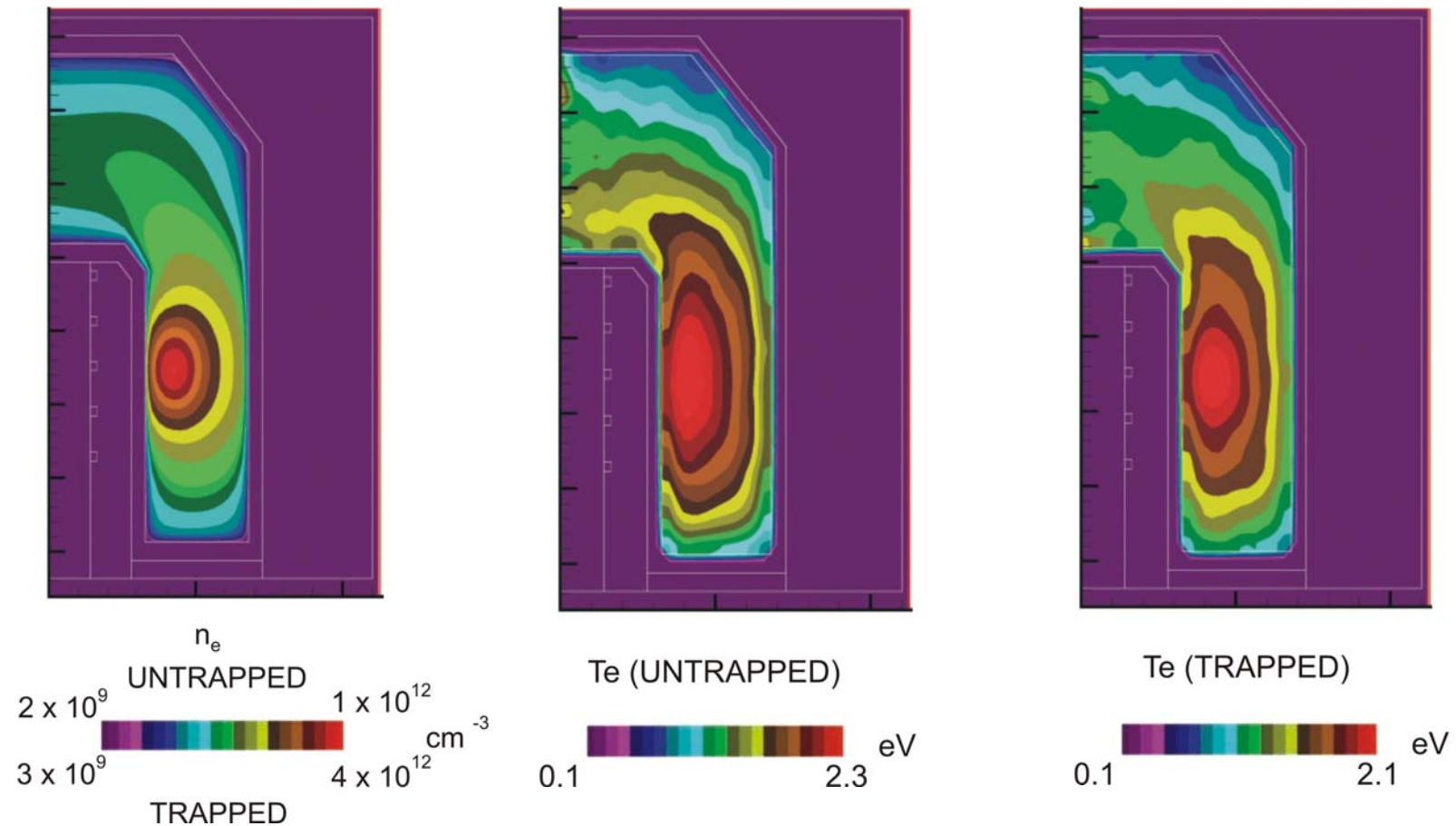
- Monte Carlo photon pseudo-particles are launched from locations proportional to Hg^* density.
- Trajectories are tracked accounting for absorption/emission based on Voigt profile.
- Null cross section techniques account for variations in absorber and perturber densities, collision frequency and gas temperature.
- Partial frequency redistribution of emitted photons.
- Isotope shifts and fine structure splitting.
- Effective lifetimes (residence times) of photons in plasma and exit spectra are calculated.
- Rate constant of radiative reaction decreased by the trapping factor (ratio of effective to natural photon lifetime for given transition).

BASE CASE CONDITIONS

- Diameter – 9 cm
- Height – 8 cm
- Initial pressure – 500 mTorr
- Initial temperature – 375 K
- Power – 50 W
- Frequency – 10 Mhz
- Initial Ar mole fraction $\cong 0.97$
- Initial Hg mole fraction (ground state) $\cong 0.03$



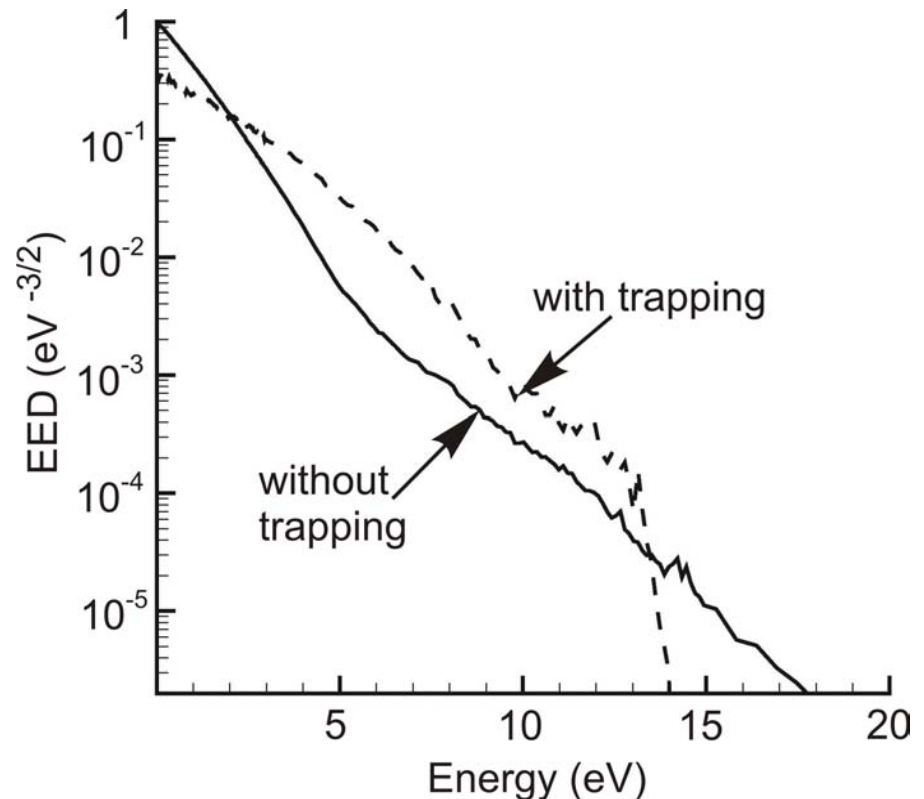
ELECTRON DENSITY AND TEMPERATURE



- The electron density goes up with trapping due to more ionization processes, while the temperature becomes more localized due to a reduced skin depth.

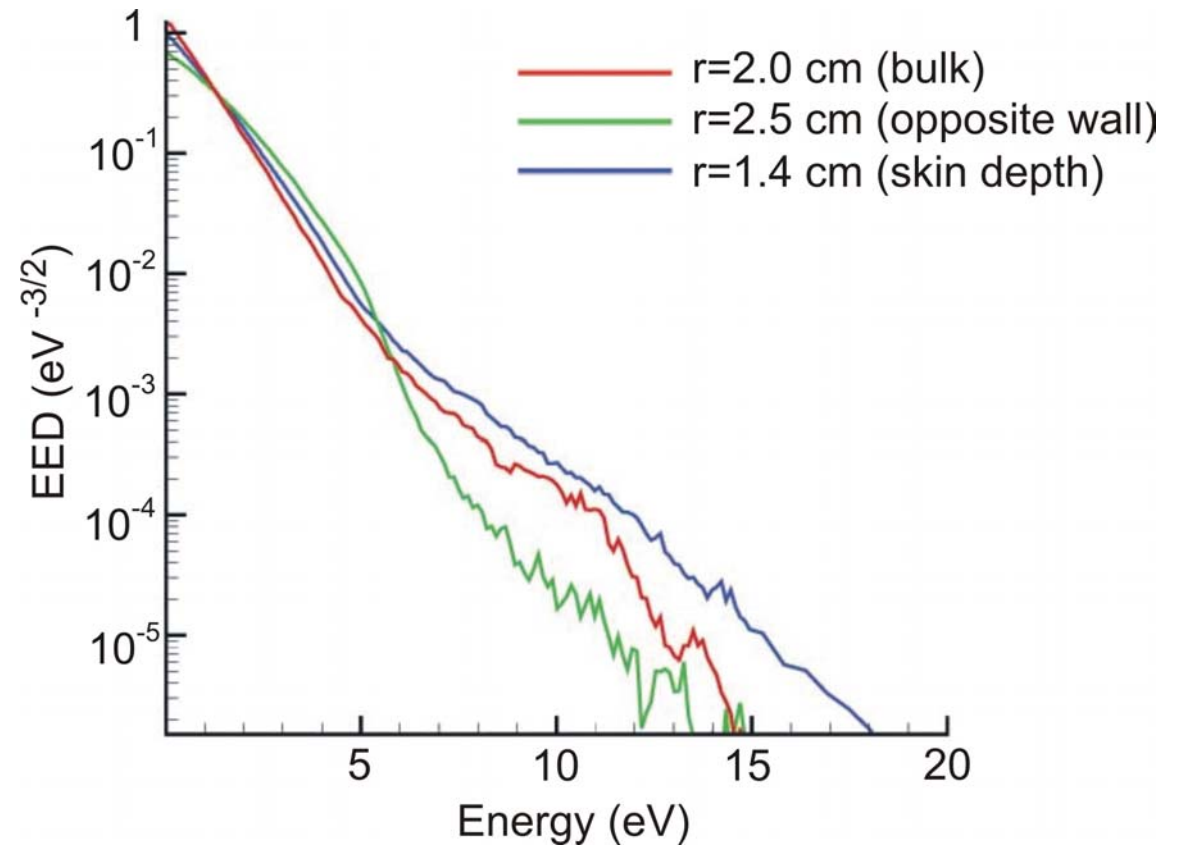
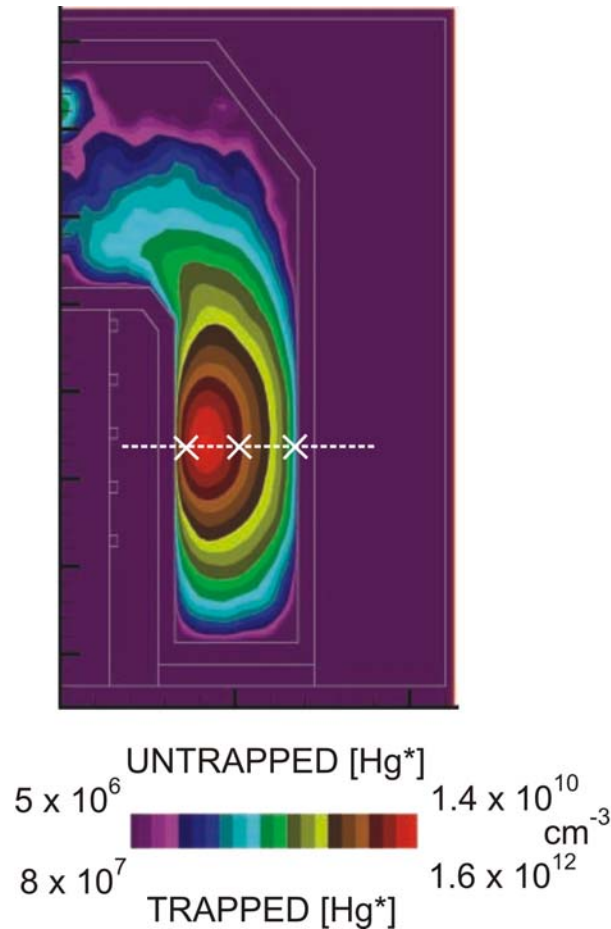
EFFECT OF RADIATION TRAPPING

- Radiation trapping leads to an increased lifetime for the $[\text{Hg}^*]$ atoms.
- There are more super-elastic collisions and hence most lower energy electrons gain an additional amount of energy ($\cong 5$ eV).



- This leads to a “bulge” in the lower energies in the EED, which is smoothed out by other inelastic collisions.
- However, since the total power deposited is a constant, higher rate of dissipation at low energies produces a decrease in the EED at higher energies.

VARIATION WITH SPATIAL POSITION



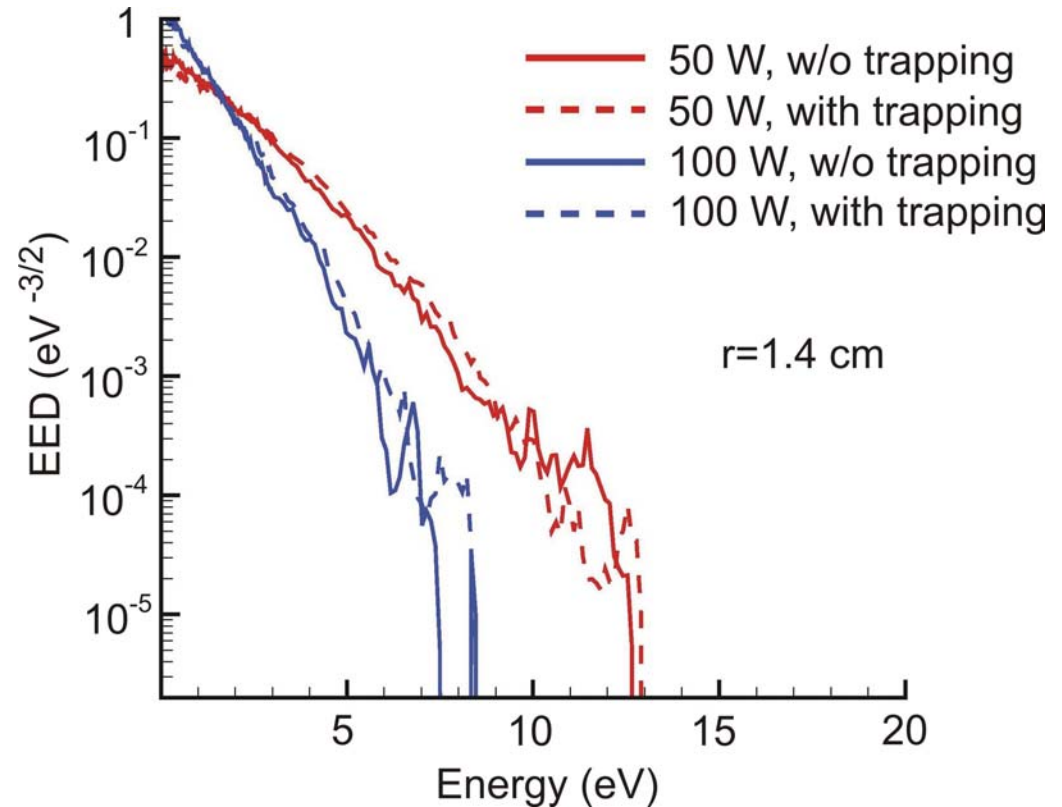
- The EED in the skin depth has a longer tail due to more stochastic heating.
- Diffusion cooling near sheath of the opposite wall depletes high energy electrons.

VARIATION WITH POWER DEPOSITED

- The EEDs shown here are for an untrapped lifetime of 125 ns ($6^3P_1 \rightarrow 6^1S_0$ transition).
- As the power increases, we have more electrons in the plasma:

$$n_e = 2.1 \times 10^{11} \text{ cm}^{-3} \text{ (50 W)}$$

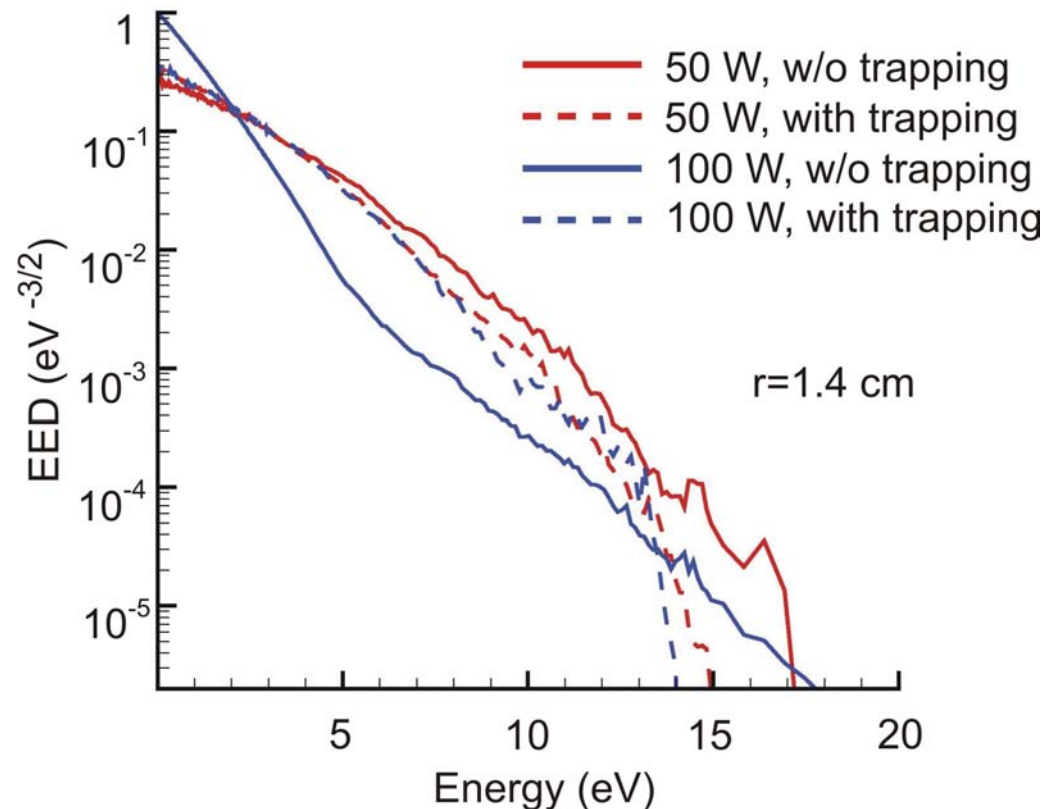
$$= 1.8 \times 10^{12} \text{ cm}^{-3} \text{ (100 W)}$$



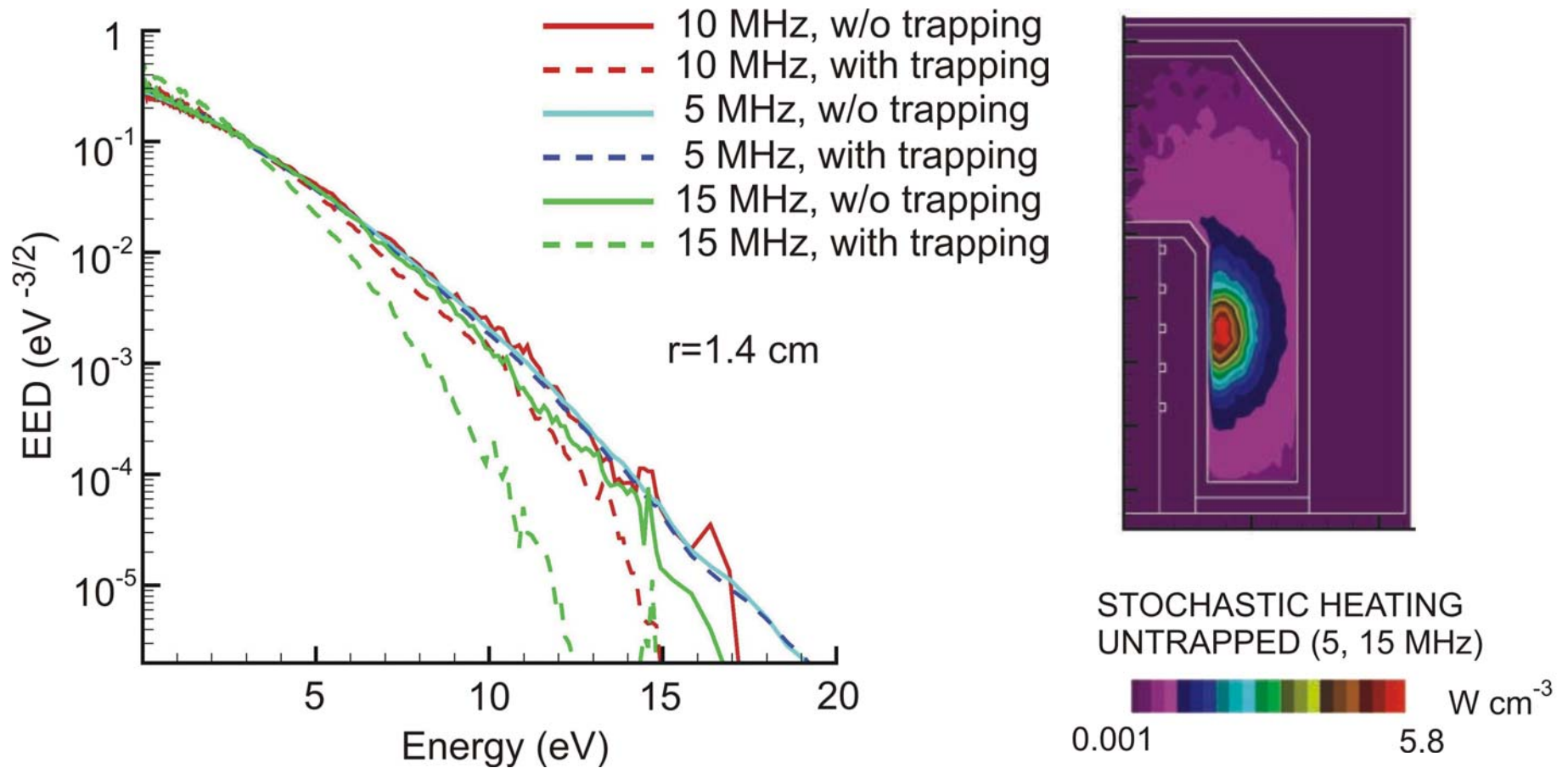
- This increased electron density leads to more Maxwellian EEDs (both trapped and untrapped) for deposition of 100 W.
- The two-temperature distribution seen at 50 W is not seen at 100 W.

VARIATION WITH POWER DEPOSITED (10 ns LIFETIME)

- The EEDs shown here are for a fictitious transition having the same energy difference as the $6^3P_1 \rightarrow 6^1S_0$ transition, but with a vacuum lifetime of 10 ns.
- The increased power deposition in the plasma leads to a doubling of electron heating in the skin depth.
- $n_e = 1.6 \times 10^{11} \text{ cm}^{-3}$ (50 W)
 $= 2.1 \times 10^{11} \text{ cm}^{-3}$ (100 W)
- This heating leads to a tail in the untrapped EED with increase in power.
- In the case of 10 ns radiative lifetime, we see that the electric field acceleration dominates the EED, while in the 125 ns case, e-e collisions determine the EED at high powers.

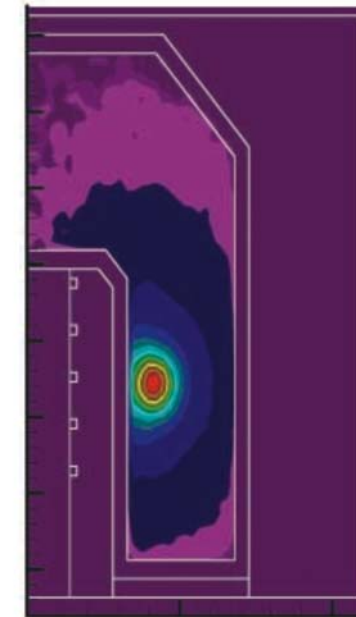
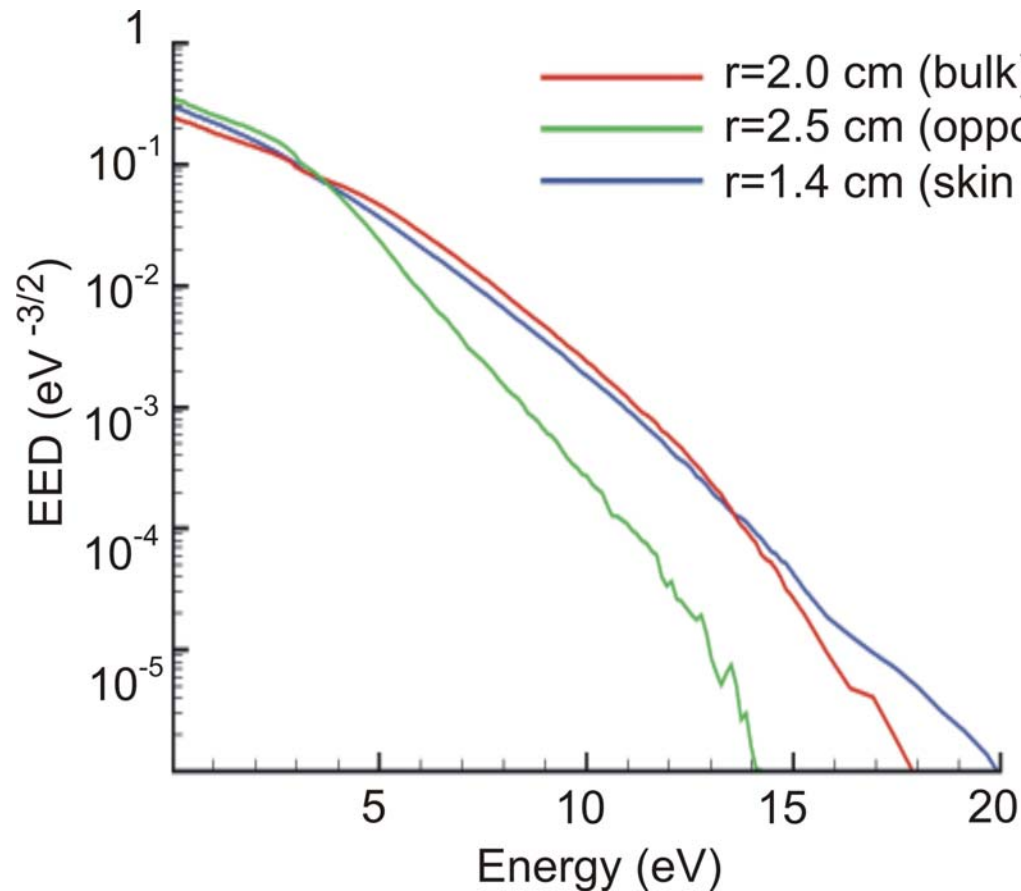


EFFECT OF ICP FREQUENCY



- With increase in frequency, the untrapped distributions do not change.
- However, the trapped EEDs are significantly different in the three cases, with lower frequencies showing a longer tail in the distribution.

EFFECT OF ICP FREQUENCY (contd.)



STOCHASTIC HEATING
TRAPPED (5 MHz)

0.001 25 W cm^{-3}

- There is more stochastic heating in the trapped EED at skin depth.
- At low frequencies, this heating offsets the reduction in the tail due to trapping.

SUMMARY

- **A Monte Carlo Resonance Radiation Transport Model has been interfaced with a plasma equipment model to model the effects of radiation trapping on EEDs.**
- **Radiation trapping is seen to affect the bulk as well as the tail of the EED due to enhancement of super-elastic collisions.**
- **The EED becomes more Maxwellian with increase in power at high radiative lifetimes. At lower lifetimes, stochastic heating dominates the tail of the EED.**
- **Increased skin depth electron heating at lower frequencies for trapped distributions significantly changes the tail of the EED.**