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**2-DIMENSIONAL MODELING OF PULSED PLASMAS  
WITH AND WITHOUT SUBSTRATE BIAS  
USING MODERATE PARALLELISM\***

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# AGENDA

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- **Introduction**
- **Description of the moderately parallel model**
- **Validation**
- **Properties of Pulsed Cl<sub>2</sub> Plasmas**
  - **With and without rf substrate bias**
  - **Bias voltage**
  - **Feedstock composition**
- **Conclusions**

# PULSED ICP PLASMAS

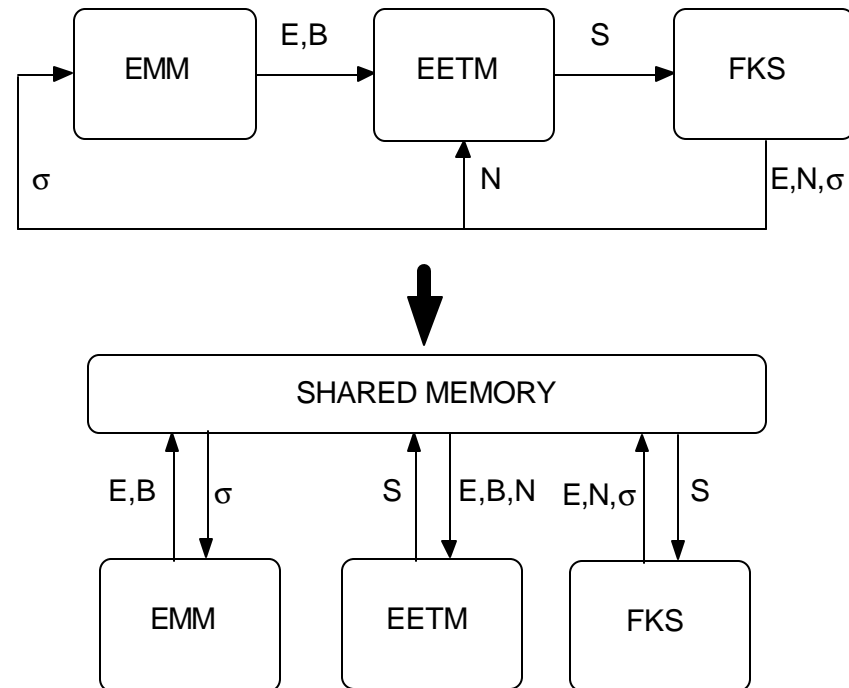
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- **Pulsed plasmas**
  - **Damage free plasma etching with better uniformity and anisotropy**
  - **Improved etch selectivity by modifying the ratio of chemical species**
  - **Additional controllable degrees of freedom : Duty cycle and modulation frequency**
  - **Reduce charge buildup on wafers and suppress notching**
  - **Reduced particle generation in the plasma**
- **Current models for investigating pulsed operation are typically global or 1-dimensional**
- **Difficult to resolve long-term transients in multi-dimensional plasma equipment models**
- **Moderately parallel algorithms for 2-dimensional hybrid models were developed to investigate long term transients.**

# DESCRIPTION OF THE PARALLEL HYBRID MODEL

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- The HPEM, a modular simulator, was parallelized by employing a shared memory programming paradigm on a Symmetric Multi-Processor (SMP) machine.
- The Electromagnetics, Electron Monte Carlo and Fluid-kinetics Modules are simultaneously executed on three processors.
- The variables updated in different modules are immediately made available through shared memory for use by other modules.
- Dynamic load balancing is implemented to equal the tasks on different processors.



# GOVERNING EQUATIONS IN HPEM

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- Continuity (heavy species) :

$$\frac{\partial N_i}{\partial t} = \nabla \cdot (N_i \vec{v}_i) + S_i$$

- Momentum (heavy species) :

$$\frac{\partial (N_i \vec{v}_i)}{\partial t} = \frac{q_i}{m_i} N_i (\vec{E}_s + \vec{v}_i \times \vec{B}_s) - \frac{1}{m_i} \nabla P_i - \nabla \cdot (N_i \vec{v}_i \vec{v}_i) - \nabla \cdot \bar{\tau}_i + \sum_j N_i N_j k_{ij} (\vec{v}_j - \vec{v}_i)$$

- Energy (heavy species) :

$$\begin{aligned} \frac{\partial}{\partial t} \left( \frac{3}{2} N_i k_B T_i \right) &= \frac{N_i q_i^2 \mathbf{n}_i}{m_i (\mathbf{n}_i^2 + \mathbf{w}^2)} E^2 + \frac{N_i q_i^2}{m_i \mathbf{n}_i} E_s^2 - \nabla \cdot \left( \frac{5}{2} k_B T_i \vec{\mathbf{j}}_i - \mathbf{I}_i \nabla T_i \right) \\ &+ \sum_i 3 \frac{m_{ij}}{m_i + m_j} N_i N_j \mathbf{n}_{ij} k_B (T_j - T_i) \pm \sum_i 3 N_i N_j R_{ij} k_B T_j \end{aligned}$$

- Drift-diffusion (electron) :

$$\frac{\partial n_e}{\partial t} = \nabla \cdot (n_e \bar{\mu}_e E_s + \bar{D}_e \nabla n_e) + S_e$$

## GOVERNING EQUATIONS IN HPEM (continued)

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- Electrons:

- Energy  $\nabla \cdot \mathbf{k} \nabla T_e + \nabla \cdot (\Gamma T_e) = P_{\text{heating}} - P_{\text{loss}}$

- Monte Carlo Simulation

$$\frac{d\vec{v}}{dt} = \frac{q_e}{m_e} (\vec{E} + \vec{v} \times \vec{B}) \quad \frac{d\vec{r}}{dt} = \vec{v} \quad v_i = \left( \frac{2\epsilon_i}{m_e} \right)^{\frac{1}{2}} \sum_{j,k} \sigma_{ijk} N_j$$

- Poisson's equation:

$$\nabla \cdot \epsilon \nabla \Phi(t + \Delta t) = \rho(t) + \frac{d\rho(t)}{t} \cdot \Delta t = \rho(t) + \sum_i q_i \Delta t [-\nabla \cdot \vec{\phi}_i + S_i]$$

- Wave equation:

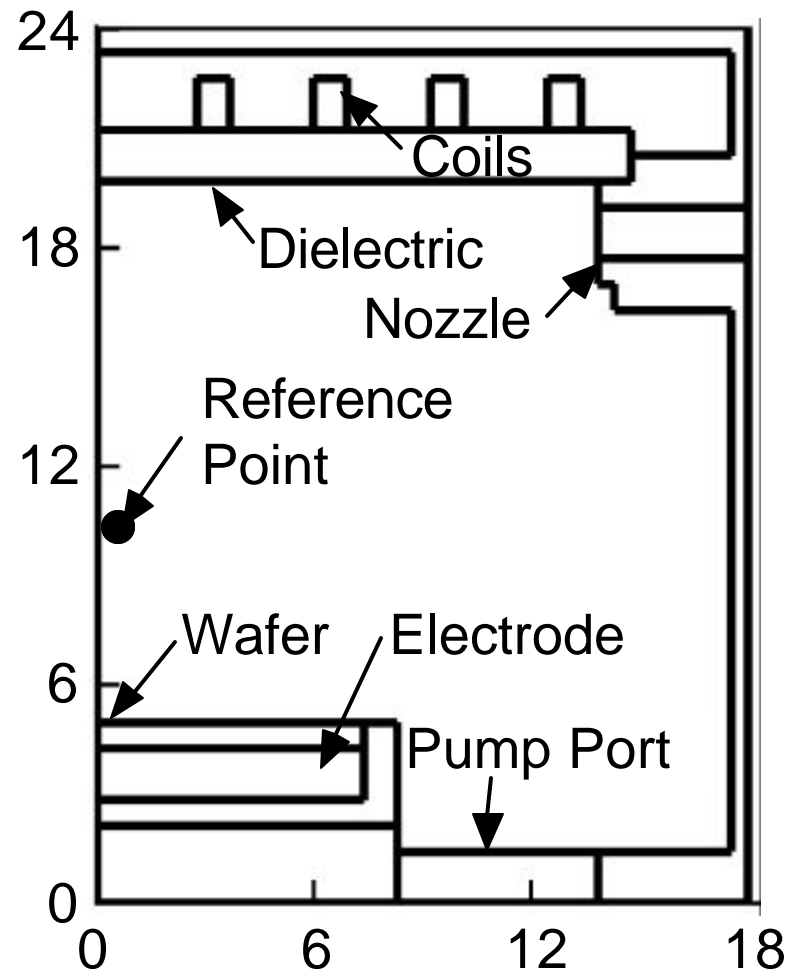
$$\nabla \cdot \frac{1}{\mathbf{m}} \nabla E = \frac{\partial^2 (\mathbf{e}_0 E)}{\partial t^2} + \frac{\partial (\overline{\mathbf{S}} E + J_0)}{\partial t}$$

- Vector potential:

$$\nabla \times \frac{1}{\mathbf{m}} \nabla \times \mathbf{A} = \mathbf{j}_B \quad \mathbf{B}_s = \nabla \times \mathbf{A}$$

# REACTOR GEOMETRY AND SIMULATION CONDITIONS

- Reactor geometry taken from Malyshev et. al.\*
- Base case conditions:
  - Peak ICP power: 600 W
  - rf bias: 250 V, 10 MHz
  - PRF: 10 kHz
  - Pressure: 10 mTorr
  - Gas flow rate: 100 sccm
  - Cl<sub>2</sub>, Ar/Cl<sub>2</sub>

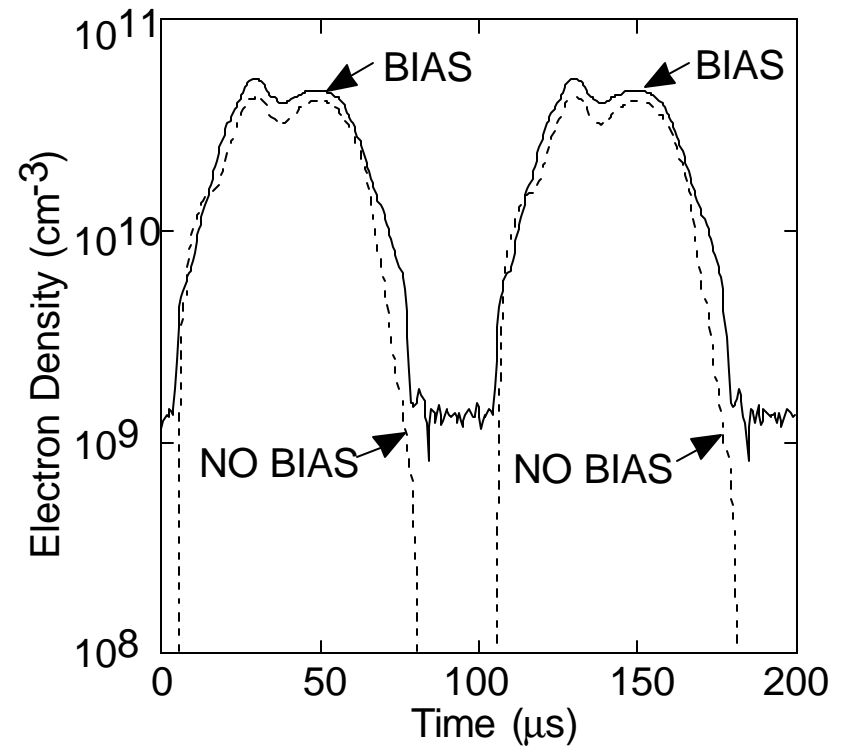


\*M. V. Malyshev, V. M. Donnelly, Plasma Sources Sci. Technol. 9, 353 (2000)

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# DYNAMICS OF PULSED PLASMAS: ELECTRON DENSITY

- ICP power at 10 MHz (rf) is pulsed at PRF of 10 kHz.
- Several pulses are required to attain periodic steady state
- $[e]$  attains steady state value in the late activeglow corresponding to cw operation
- $[e]$  decays several orders of magnitude, as the electrons are lost due to dissociative attachment
- With substrate bias,  $[e]$  attains a steady state value in the late after glow corresponding to capacitive mode



- $\text{Cl}_2$ , 10 mTorr, 600 W, 10 MHz, 10 kHz/50%

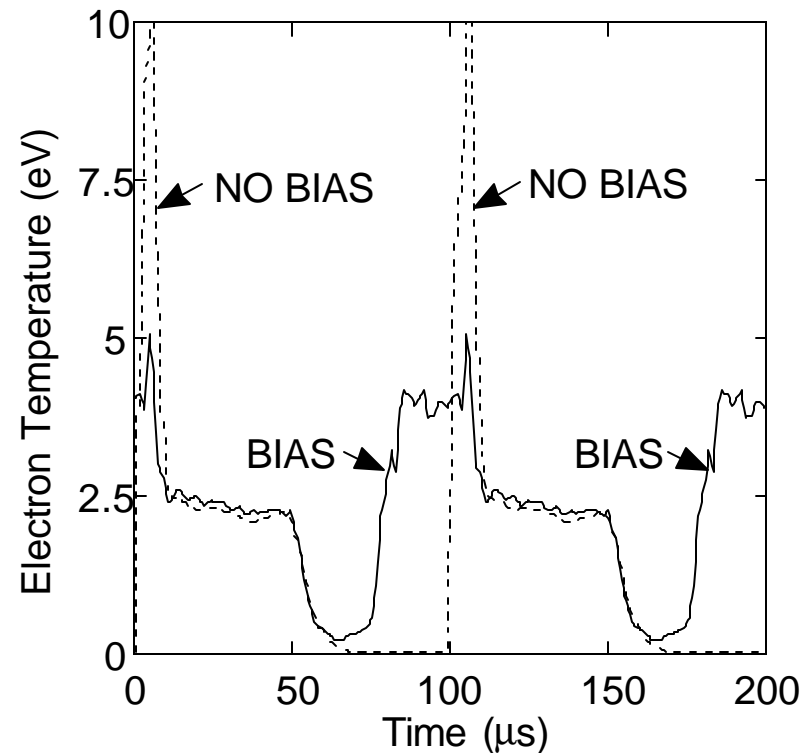
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# DYNAMICS OF PULSED PLASMAS: ELECTRON TEMPERATURE

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- $T_e$  peaks at the leading edge, as the power deposition occurs into a smaller inventory of electrons
- No substrate bias results in higher peak  $T_e$ , as the electrons decay away several orders of magnitude lower in the late afterglow
- $T_e$  is similar during late activeglow and early afterglow period
- With substrate bias,  $T_e$  increases in the late afterglow due to sheath heating
- Electrons thermalize to gas temperature without substrate bias



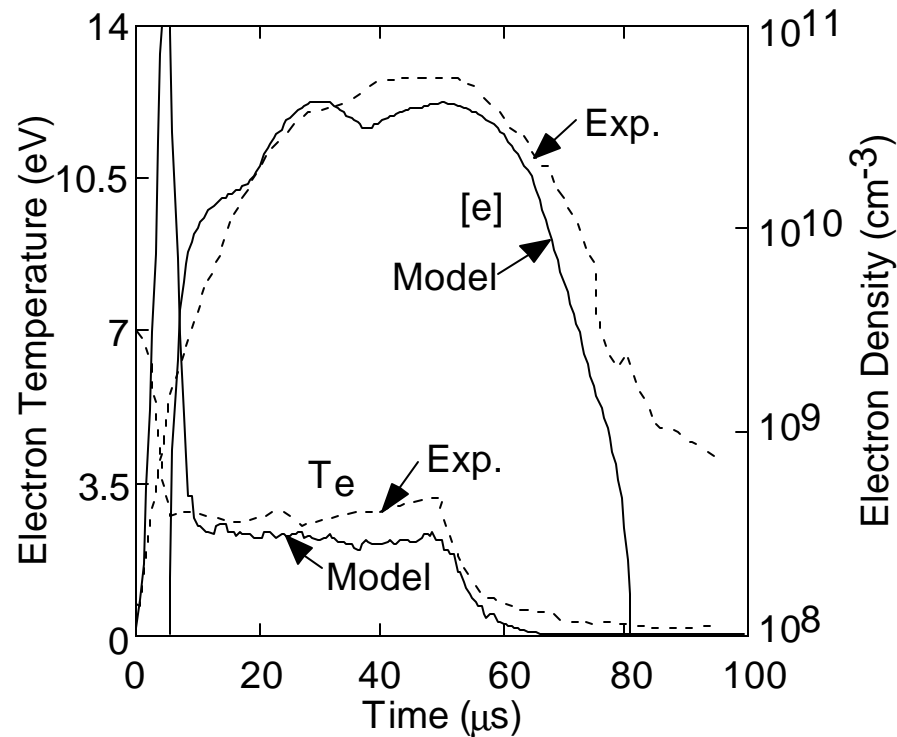
- Electron Temperature

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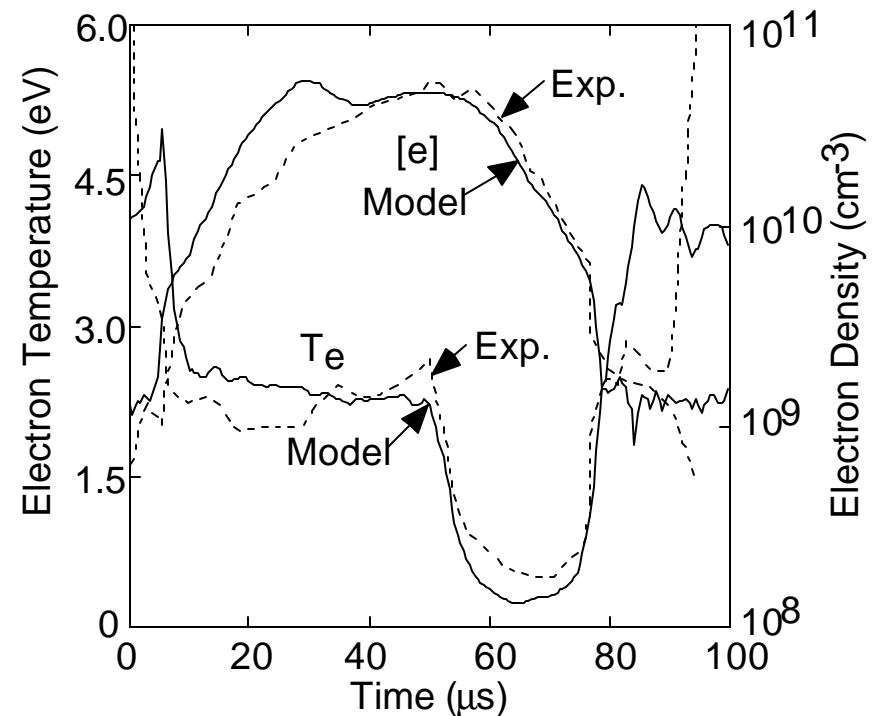
# VALIDATION OF THE MODEL: PHYSICS

- Model results compare well with experiments\*
- With substrate bias,  $T_e$  rises to values above the steady state value



- Without rf bias

- Cl<sub>2</sub>, 10 mTorr, 600 W, 10 MHz, 10 kHz/50%



- With rf bias (250 V, 10 MHz)

\* M. V. Malyshev, V. M. Donnelly, Plasma Sources Sci. Technol. 9, 353 (2000)

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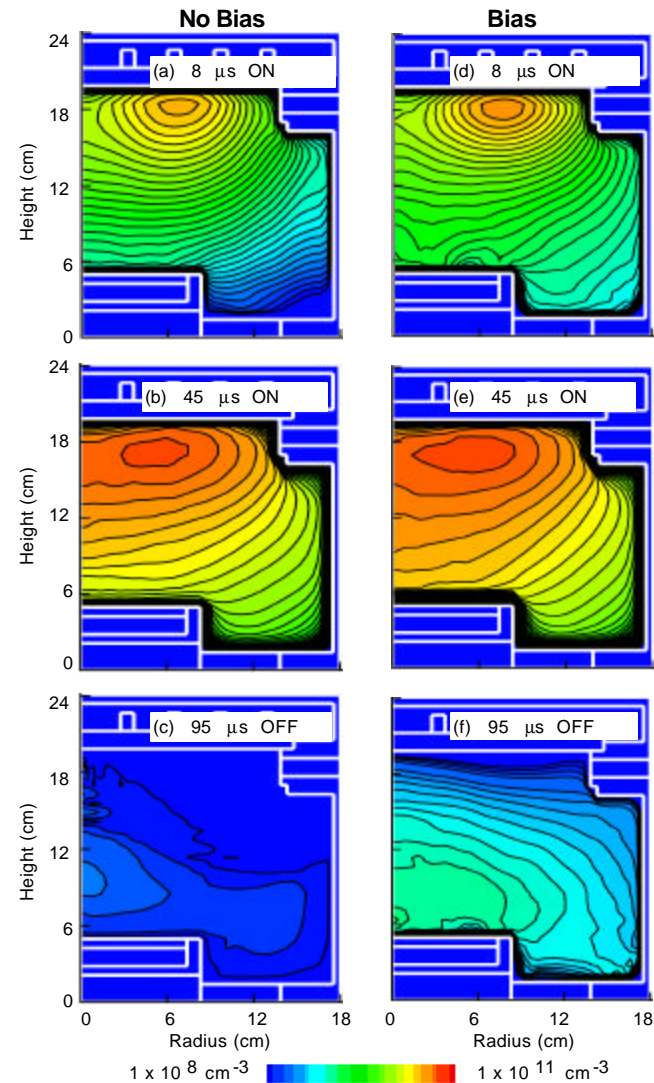
# SHEATH HEATING IN PULSED PLASMAS WITH SUBSTRATE BIAS

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- Sheath heating scales with the square of the sheath speed,  $v_s$
- Sheath thickness  $\mathbf{l}$  scales as  $n_e^{-1/2}$
- Sheath speed  $v_s \gg \mathbf{wl}$
- Total sheath heating  $H$ , therefore scales as  $H \sim v_s^2 n_e$  and is not a function of electron density
- Specific heating rate (per electron),  $h$ , scales as  $H/n_e \sim 1/n_e$
- As the electron density decays, primarily by dissociative attachment to  $\text{Cl}_2$ , the sheath thickness  $\mathbf{l}$  increases
- Oscillating sheath produces a net increase in specific heating rate  $h$  and hence an increase in (or slowing in rate of decrease in)  $T_e$

# Cl<sub>2</sub> PULSED PLASMAS WITH AND WITHOUT RF BIAS: ELECTRON DENSITY

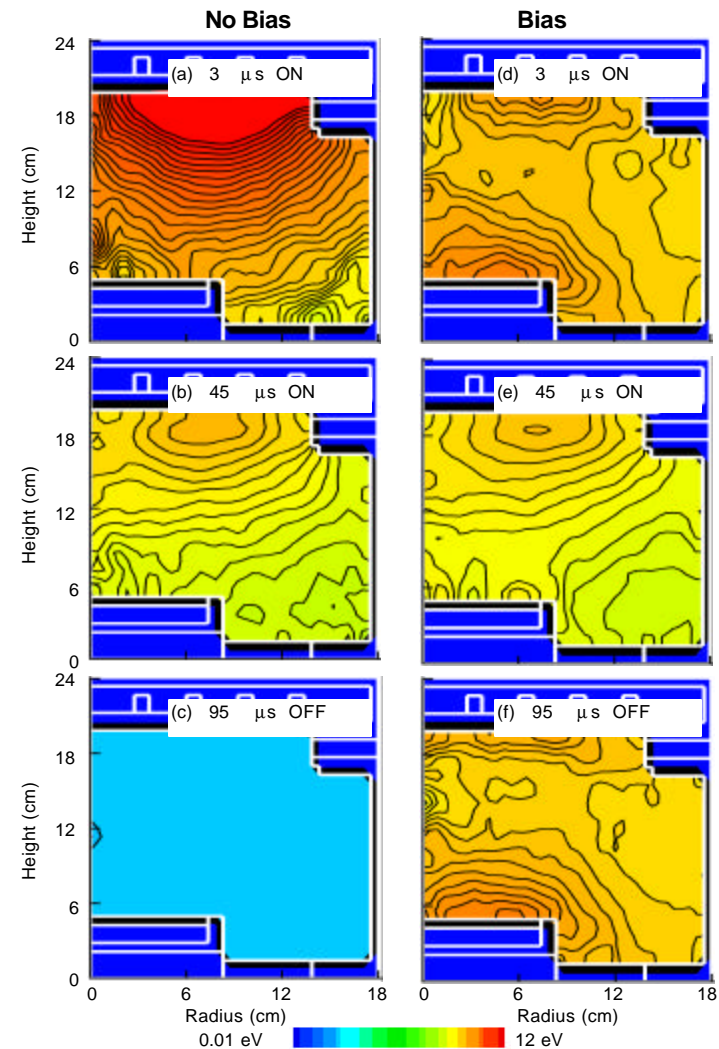
- The peak [e] migrates to below the coils during "power-on" where the source is maximum.
- [e] is similar with and without substrate bias during the activeglow and early afterglow phase.
- As the power is turned off, in the early afterglow ambipolar losses dominate over generation of electrons.
- In the late afterglow, sheath heating dominates and the plasma transitions from inductively coupled to capacitively coupled mode with substrate bias.



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# Cl<sub>2</sub> PULSED PLASMAS WITH AND WITHOUT RF BIAS: ELECTRON TEMPERATURE

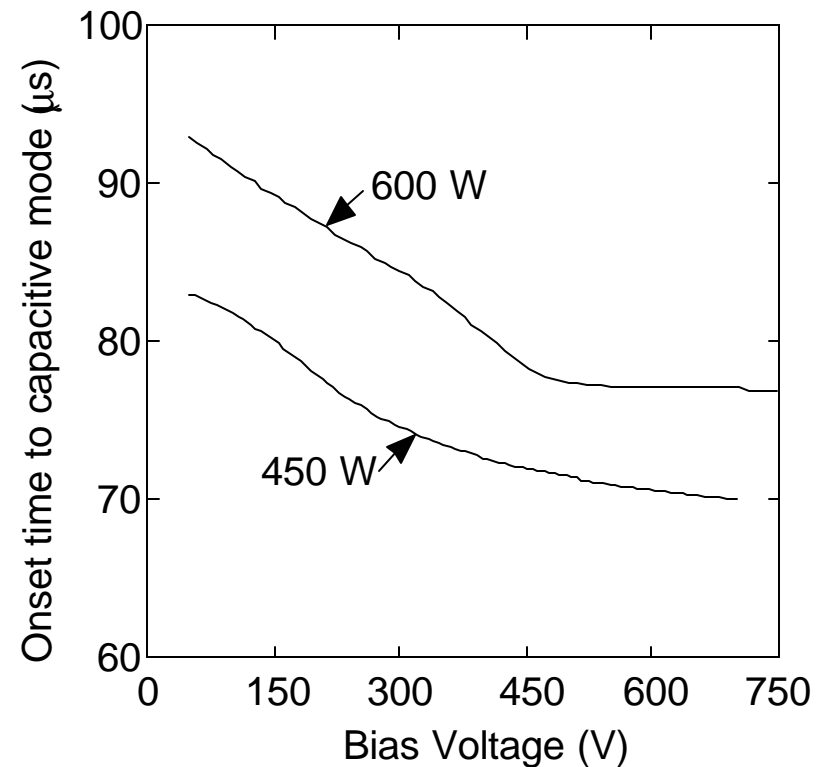
- The peak in  $T_e$  at the leading edge is due to power deposition into smaller inventory of electrons.
- Sheath heating prevails, even after 10  $\mu\text{s}$  into the power-on period, owing to low  $n_e$ .
- At 45  $\mu\text{s}$ , the  $T_e$  profile looks similar to no substrate bias case
- After 25  $\mu\text{s}$  into the power-off period, sheath heating begins
- $T_e$  increases in the late afterglow, as power is deposited into electrons near the substrate by oscillating sheath.



# EFFECT OF BIAS VOLTAGE AND ICP POWER ON ONSET TIME

- As the bias voltage is increased, the onset time to the capacitive coupling mode,  $t_c$ , decreases.
- This is attributed to the greater sheath heating at higher bias voltages.
- As the ICP power is varied from 450 W to 600 W,  $t_c$  increases as the peak electron density in the late active glow is higher (thinner sheath).
- At higher ICP powers, the plasma is more dissociating and less attaching, which increases the  $[e]$  in the late active glow and increases  $t_c$ .

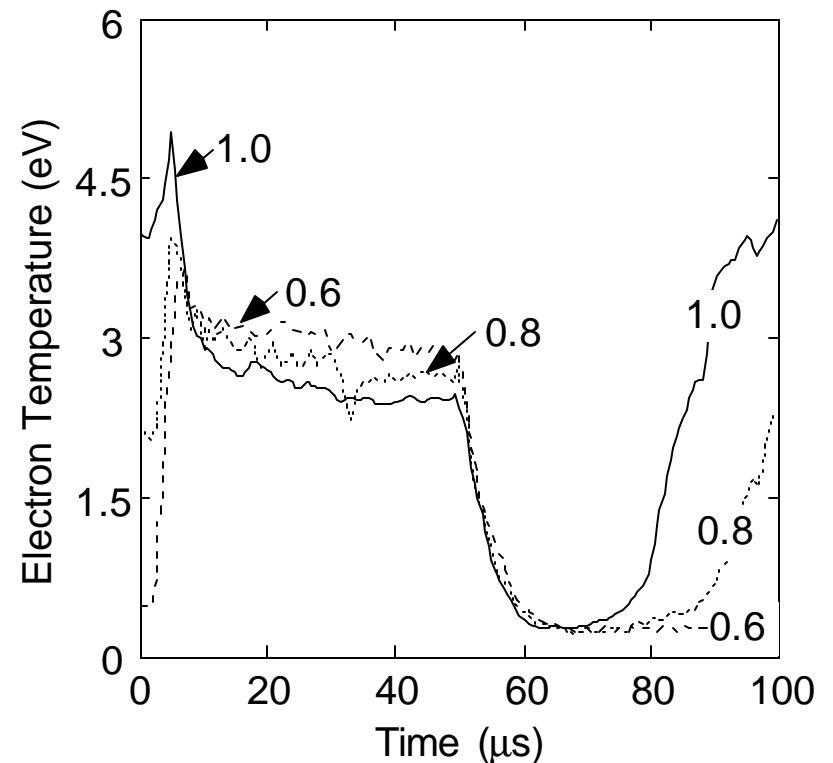
- $\text{Cl}_2$  pulsed plasma
- 10 mTorr, 100 sccm
- PRF: 10 kHz
- Duty cycle: 50%
- Bias: 250 V



# EFFECT OF GAS MIXTURES ON ONSET TIME

- As the Ar fraction is increased,  $t_c$  increases significantly.
- For Ar/Cl<sub>2</sub> =40/60,  $t_c$  is longer than the pulse off time, due to the higher [e] in afterglow (thinner sheath).
- As sheath heating is reduced, negative ions can be extracted during the pulse off period.
- Bias voltage, frequency and feedstock composition need to be optimized to investigate the possibility of negative ion extraction with substrate bias.

- Ar/Cl<sub>2</sub> pulsed plasma
- 450 W, 10 mTorr
- 100 sccm
- PRF: 10 kHz
- Duty cycle: 50%
- Bias: 250 V



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# CONCLUSIONS

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- A new 2-D hybrid model was developed to address transients based on moderate computational parallelism.
- Computational studies were performed for pulsed operation of  $\text{Cl}_2$  and  $\text{Ar}/\text{Cl}_2$  ICPs.
- $T_e$  at the leading edge is nearly twice the steady state value.
- In electronegative plasmas, electron-ion plasma in the activeglow becomes ion-ion plasma in the afterglow.
- For pulsed  $\text{Cl}_2$  plasmas with continuous substrate bias, sheath heating was observed to be predominant in the late afterglow.
- The  $t_c$  increases with ICP power as the plasma is more dissociated in the late activeglow and it takes more time for electron density to decay away.
- The extraction of negative ions in the afterglow is difficult with a continuous substrate bias as the sheaths typically do not collapse.